

## Article

# The Equivalent Effect of Initial Condition Coupling on the Laminar Burning Velocity of Natural Gas Diluted by CO<sub>2</sub>

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**Abstract:** Initial temperature has a promoting effect on laminar burning velocity, while initial pressure and dilution rate have an inhibitory effect on laminar burning velocity. Equal laminar burning velocities can be obtained by initial condition coupling with different temperatures, pressures and dilution rates. This paper analysed the equivalent distribution pattern of laminar burning velocity and the variation pattern of an equal weight curve using the coupling effect of the initial pressure (0.1–0.3 MPa), initial temperature (323–423 K) and dilution rate (0–16%). The results show that, as the initial temperature increases, the initial pressure decreases and the dilution rate decreases, the rate of change in laminar burning velocity increases. The equivalent effect of initial condition coupling can obtain equal laminar burning velocity with an dilution rate increase (or decrease) of 2% and an initial temperature increase (or decrease) of 29 K. Moreover, the increase in equivalence ratio leads to the rate of change in laminar burning velocity first increasing and then decreasing, while the increases in dilution rate and initial pressure make the rate of change in laminar burning velocity gradually decrease and the increase in initial temperature makes the rate of change in laminar burning velocity gradually increase. The area of the region, where the initial temperature influence weight is larger, gradually decreases as the dilution rate increases, and the rate of decrease gradually decreases.

**Keywords:** natural gas; laminar burning velocity; equivalent effect; influence weight



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## 1. Introduction

Natural gas is considered the oil alternative fuel with the most potential due to its advantages of eco-friendliness, cleanliness, low carbon emission, high efficiency and low price, which have attracted the attention of scholars at home and abroad [1–3]. The laminar burning characteristics of natural gas are of great significance for understanding the inherent physical and chemical properties, the flame propagation process and the chemical kinetics [4–6]. Domestic and international scholars have conducted a lot of research regarding the influence of initial pressure ( $P_u$ ), initial temperature ( $T_u$ ), mixed gas composition and concentration on laminar burning velocity ( $u_L$ ). Han et al. [7] analysed the effect of initial temperature (323–423 K) on the laminar flame of premixed natural gas through experimental studies. The results show that, with the increase in  $T_u$ ,  $u_L$  gradually increases. Hermanns et al. [8] summarised the available measurements of laminar burning velocities in CH<sub>4</sub> + H<sub>2</sub> + O<sub>2</sub> + N<sub>2</sub> flames at a temperature range of 298–418 K performed using a heating flux method. The results show that the increase in  $T_u$  increases  $u_L$  under different equivalence ratios. Halter et al. [9] analysed the effect of  $P_u$  (0.1–0.5 MPa) on the laminar flame of CH<sub>4</sub>/air mixtures through experimental studies. The results show that, with the increase in  $P_u$ ,  $u_L$  gradually decreases. Xie et al. [10] carried

out a chemical kinetic modelling study of laminar burning characteristics for CH<sub>4</sub>/CO<sub>2</sub> mixtures at elevated pressure by CHEMKIN coupling with a detailed chemical reaction mechanism. The results show that  $u_L$  decreases with increasing pressure under high pressure. Chan et al. [11] studied the effects of CO<sub>2</sub> diluent on  $u_L$  of CH<sub>4</sub>/air premixed flames utilizing experimentation and kinetic modelling. The results show that  $u_L$  of the methane and air mixture decreases as the CO<sub>2</sub> dilution rate (DR) increases. Zhou et al. [12] conducted a study on the effect of diluents (N<sub>2</sub>/CO<sub>2</sub>) on the laminar flame speed of a H<sub>2</sub>/CO/CH<sub>4</sub>/air premixed flame using an outwardly propagating spherical flame and the CHEMKIN package. The results show that laminar flame speed decreases with the increase in N<sub>2</sub>/CO<sub>2</sub> dilution ratios and that CO<sub>2</sub> dilution has a stronger dilution effect, thermal effect and chemical effect than those of N<sub>2</sub> dilution. Huang et al. [13] studied the laminar flame characteristics of natural gas–air flames in a constant-volume bomb at normal temperature and pressure. The results show that  $u_L$  tends to increase first and then decrease with the increase in the equivalence ratio ( $\Phi$ ), and the maximum value is obtained between  $\Phi$  from 1.0 to 1.1. Dirrenberger et al. [14] presented new experimental measurements of the laminar flame velocity of natural gas with equivalence ratios from 0.6 to 2.1 performed by the heat flux method. The results show that, with the increase in  $\Phi$ ,  $u_L$  increases first and then decreases. This pattern maintains good consistency with other test data from the literature.

According to the above review,  $T_u$  has a promoting effect on  $u_L$  while  $P_u$  and DR have an inhibitory effect on  $u_L$ . Equal laminar burning velocities can be obtained by initial condition coupling with different temperatures, pressures and DRs. However, the literature lacks corresponding research results on the equivalent effect on laminar burning. There are much domestic and foreign research on the influence of parameters such as  $T_u$ ,  $P_u$  and DR on combustion, but most of the research focuses on the influence of a single parameter on laminar burning, and it is difficult to obtain a quantitative equivalent relationship. Based on the equivalent laminar burning concept, this paper analysed the equivalent distribution pattern of  $u_L$  and the variation pattern of the equal weight curve by the coupling effects of  $P_u$  (0.1–0.3 MPa),  $T_u$  (323–423 K) and DR (0–16%). Relevant data support and engineering reference are provided for revealing the influence of the coupling mechanism of initial parameters on the laminar burning process, which is of great significance.

## 2. Experimental Setup

Figure 1 is a schematic of the test system, which is mainly composed of a constant volume chamber (CVC), a temperature monitoring system, an ignition system, a data acquisition system and a Schlieren imaging system [15]. The temperature monitoring system includes a K-Type thermocouple and a proportional-integral-derivative (PID) temperature controller, which is capable of maintaining the  $T_u$  error within  $\pm 3$  k. The parameters of the ignition system are as follows: the ignition voltage is 14 V, provided by a stabilized power supply; the ignition pulse width is 3 ms; the ignition electrode diameter is 2 mm; and the gap between the ignition electrodes is 3 mm. The data acquisition system includes a pressure sensor (KISTLER 6125C), a data acquisition card (Data acquisition, NiUSB-6365, sampling frequency of 100,000 Hz) and a charge amplifier (KISTLER 5018A). The Schlieren imaging system includes an illuminator (100 W Power), two concave mirrors (Focal length 110 mm), two plane mirrors and a high-speed digital camera (Phantom V7.3, 10,000 fps, resolution 512  $\times$  512 pixels). Table 1 shows the main parameters of the CVC.

According to the above review,  $T_u$  has a promoting effect on  $u_L$  while  $P_u$  and  $DR$  have an inhibitory effect on  $u_L$ . Equal laminar burning velocities can be obtained by initial condition coupling with different temperatures, pressures and  $DR$ s. However, the literature lacks corresponding research results on the equivalent effect on laminar burning. There are much domestic and foreign research on the influence of parameters such as  $T_u$ ,  $P_u$  and  $DR$  on combustion, but most of the research focuses on the influence of a single parameter on laminar burning, and it is difficult to obtain a quantitative equivalent

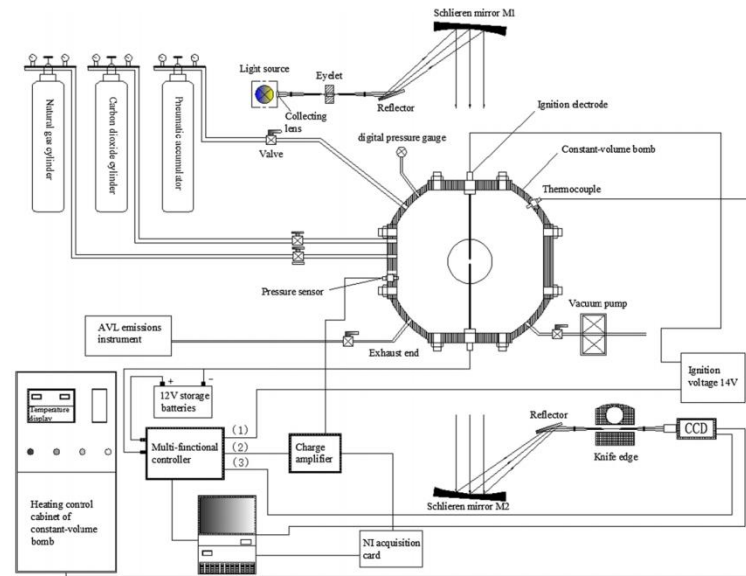


Figure 1. The experimental setup.

Table 1. Basic parameters of the constant volume chamber (CVC).

Parameter (Unit)	Value
Inner diameter (mm)	350
Volume (L)	22.422.4
Maximum heating temperature (K)	600
Maximum pressure (MPa)	4
Effective diameter of windows (mm)	$\Phi 120$

The experiment uses a mixture of  $CO_2$ , natural gas and compressed air to carry out premixed combustion research in order to meet the needs of modern society for natural gas engine performance simulation.  $CO_2$  is used as an inert gas to reduce the oxygen concentration of the reactant, mainly to simulate the exhaust gas recirculation (EGR) technology of the engine. Compared with  $N_2$ ,  $CO_2$  has a greater impact on the laminar combustion of mixed gas, which is closer to the actual use of real EGR technology. In the test, gaseous  $CO_2$ , natural gas and compressed air are charged to the CVC to the specified pressure sequentially according to the law of partial pressure. The chemical reaction formula of the reaction between natural gas ( $CH_4$ ) and oxygen ( $O_2$ ) is as follows:



From the metering ratio of this reaction formula, it can be seen that 1 mol methane needs to consume 2 mol oxygen gas and that the source of oxygen is air. After calculation, at  $\Phi$  of 1, the complete oxidation of 1 mol of methane needs to consume 9.524 mol of air.

The partial pressure of carbon dioxide is expressed as follows:

$$P_{CO_2} = DR \times P_u \tag{2}$$

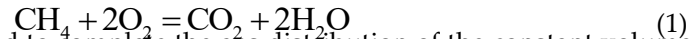
When the equivalence ratio is  $\Phi$ , the oxidation of 1 mol methane needs to consume  $9.524/\Phi$  mol air, so the partial pressure of methane when the equivalence ratio is  $\Phi$  is as follows:

$$p_{CH_4} = \frac{P_u - P_{CO_2}}{1 + 9.524/\phi} \tag{3}$$

The compressed air partial pressure is obtained from the total initial pressure minus other gas partial pressures:

$$P_{air} = P_u - P_{CO_2} - P_{CH_4} \tag{4}$$

centration of the reactant, mainly to simulate the exhaust gas recirculation (EGR) technology of the engine. Compared with N<sub>2</sub>, CO<sub>2</sub> has a greater impact on the laminar combustion of mixed gas, which is closer to the actual use of real EGR technology. In the test, gaseous CO<sub>2</sub>, natural gas and compressed air are charged to the CVC to the specified pressure sequentially according to the law of partial pressure. The chemical reaction formula of the reaction between natural gas (CH<sub>4</sub>) and oxygen (O<sub>2</sub>) is as follows:



The above method is used to complete the gas distribution of the constant volume chamber (CVC). Meanwhile, the gas is gradually heated to T<sub>0</sub> at that time, and methane needs to consume 2 mol oxygen gas and the combustion is completed. After the calculation, gas is discharged for "scrubbing" and returned more than three times to ensure the chamber cleanliness [16]. To reduce the test error [17–19], the radius of this test is 6–25 mm. The partial pressure of carbon dioxide is expressed as follows:

$$P_{\text{CO}_2} = DR \times Pu \tag{2}$$

### 3. Data Processing

#### 3.1. Extraction of Flame Radius

When the equivalence ratio is  $\Phi$ , the oxidation of 1 mol methane needs to consume 9.524/ $\Phi$  mol air, so the partial pressure of methane when the equivalence ratio is  $\Phi$  is as follows:

Figure 2 illustrates the process of obtaining a Schlieren image of the propagation flame radius using the commercial mathematical software MATLAB. This paper chose the Canny operator because of its high precision [20]. Before detecting the image boundary, five steps were applied: background removal, greyscale processing (the threshold value is 10), flame front extraction, boundary identification and fitting [21]. In the radius calculation, the horizontal line was rotated clockwise by 0 degrees, 60 degrees and 120 degrees to obtain 3 diameters (6 radii), and the instantaneous flame radius  $R_u$  was obtained by averaging the 6 radius values.

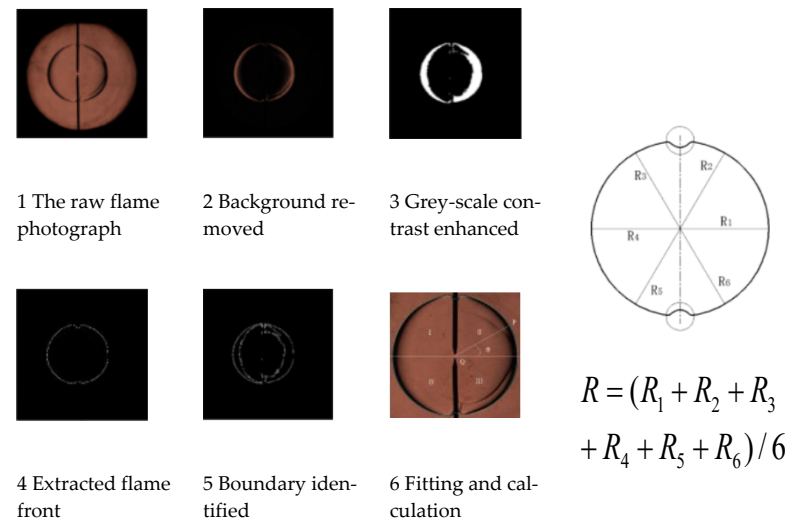


Figure 2. Diagram of flame radius extracting data calculating.

#### 3.2. Data Calculation

In the spherical diffusion flame, the propagation rate of the tensile flame is given [22], shown below:

$$S_n = dR_u / dt \tag{5}$$

where  $t$  is the time.

$$S_n = dR_u / dt \tag{4}$$

For the outwardly propagating spherical flame, the flame stretch rate can be simplified as follows [23]:

For the outwardly propagating spherical flame, the flame stretch rate can be simplified as follows [23]:

in which  $A$  is an infinitely small area on the flame and  $\kappa = 2/R_u$  is the curvature of the flame front.

To obtain the unstretched flame propagation velocity and the Markstein length, according to the literature [24], use the classical formula, shown as follow:

$$S_1 - S_n = L_b K \tag{7}$$

However, there is a certain theoretical error in this method, so the recommendation of Chen [25] is adopted. For most cases of the mixture with Lewis number  $Le < 1$  or close to 1, the nonlinear method proposed by Kelley et al. [26] is as follows:

$$\ln(S_n) = \ln(S_1) - S_1 L_b \kappa / S_n \quad (8)$$

For most cases of the mixture with  $Le > 1$ , another nonlinear formula [27] is used, shown as follow:

$$S_n = S_1 - S_1 L_b \kappa \quad (9)$$

where  $Le$  is as follows:

$$Le = \lambda / \rho_u c_p D_m = D_T / D_m \quad (10)$$

The un-stretched  $u_L$  can be calculated as follow [28]:

$$u_L = S_1 (\rho_b / \rho_u) = S_1 / \sigma \quad (11)$$

To further evaluate the influence of the coupling relationship of  $P_u$ ,  $T_u$  and  $DR$  on  $u_L$ , this paper introduces the variation of laminar burning velocity  $\Delta u_{s1-s2}$ , defined as follow:

$$\Delta u_{s1-s2} = u_{s1} - u_{s2} \quad (12)$$

where  $u_{s1}$  and  $u_{s2}$  are the  $u_L$  of  $s1$  and  $s2$ , respectively, and  $\Delta u_{s1-s2}$  is the variation of  $u_L$  between  $u_{s1}$  and  $u_{s2}$ .

To evaluate the changes in  $T_u$ ,  $P_u$  and  $DR$  under the equivalent  $\Delta u_{s1-s2}$ , define equations as follow:

$$\Delta T_{s1-s2} = |T_{s1} - T_{s2}| \quad (13)$$

$$\Delta P_{s1-s2} = |P_{s1} - P_{s2}| \quad (14)$$

$$\Delta DR_{s1-s2} = |DR_{s1} - DR_{s2}| \quad (15)$$

where  $T_{s1}$ ,  $T_{s2}$ ,  $P_{s1}$ ,  $P_{s2}$ ,  $DR_{s1}$  and  $DR_{s2}$  represent the initial temperatures, initial pressures and dilution rates for the  $u_L$  of  $s1$  and  $s2$ , respectively, and  $\Delta T_{s1-s2}$ ,  $\Delta P_{s1-s2}$  and  $\Delta DR_{s1-s2}$  are the corresponding  $T_u$ ,  $P_u$  and  $DR$  of the variation in  $u_L$  between  $u_{s1}$  and  $u_{s2}$ .

To further analyse the influence of the coupling relationship of  $P_u$ ,  $T_u$  and  $DR$  on the  $u_L$ , define the equations as follow:

$$RT_{s1-s2} = \frac{\Delta u_{s1-s2}}{\Delta T_{s1-s2}} \quad (16)$$

$$RP_{s1-s2} = \frac{\Delta u_{s1-s2}}{\Delta P_{s1-s2}} \quad (17)$$

$$RDR_{s1-s2} = \frac{\Delta u_{s1-s2}}{\Delta DR_{s1-s2}} \quad (18)$$

where,  $RT_{s1-s2}$ ,  $RP_{s1-s2}$  and  $RDR_{s1-s2}$  represent the variations in  $u_L$  between  $u_{s1}$  and  $u_{s2}$  per unit temperature, unit pressure and unit dilution rate, respectively.

### 3.3. Chemical Kinetic Model

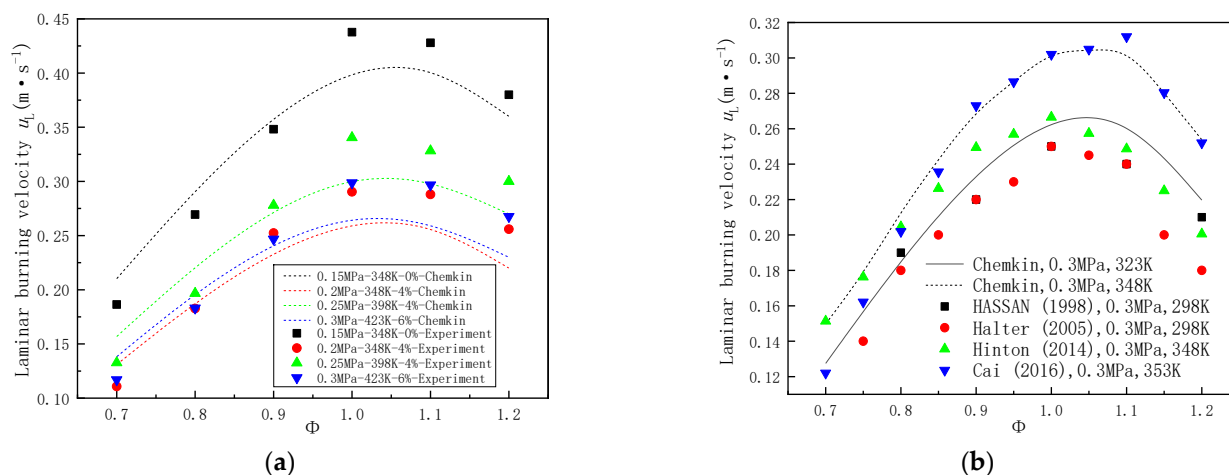
Chemkin (GRI\_mech 3.0) was applied in this study. GRI\_mech is a series of mechanisms aiming at combustion of methane that were proposed by Gas Research Institute, and GRI\_mech 3.0 is the latest version. GRI\_mech 3.0 mechanism contains 53 components and 325 elementary reactions and works well in the combustion of methane, carbon monoxide, hydrogen, etc.

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#### 4. Results and Discussion

##### Analysis of Laminar Burning Velocities Distribution of Natural Gas

Figure 3a shows the data comparison between the test and simulation. The results exhibit very good consistency between the simulated and measured values of  $u_L$ . It can be seen from the figure that the maximum error between the simulation model and the test results is 5%, within a reliable range [29–31]. To further improve the accuracy of the simulation, Figure 3b shows the deviation of measured values by different groups from that predicted by the simulation based on GRI\_mech 3.0. The results show that the difference between the simulated data (0.3 MPa and 323 K) and HASSAN's [32] and Halter's [9] (0.3 MPa and 298 K) is found to be within 4.7 cm/s at a temperature difference of 25 K. Good agreement is found between our data and Cai's [33] measurements except for the  $\Phi$  of 0.7, where our  $u_L$  is 2.9 cm/s (error 9%) higher, as the temperature difference is 5 K. The deviation in numerical data is only slightly larger than that of Hinton's [34], especially when the equivalent is relatively high, but the maximum difference is only 5 cm/s. It can be seen that the simulation model maintains good consistency with the experimental results and literature data and can be used for analysis under certain working conditions.



**Figure 3.** Validation of present simulation studies. (a) compared with experimental data and (b) compared with previous data.

Figure 4 shows the equivalent distribution pattern of  $u_L$  under the coupling effect of  $T_u$  and  $P_u$  when the equivalence ratios are 0.9, 1.0, and 1.1 and  $\Phi$  is 0%. As shown in the figure, the equivalent ratios are 0.9, 1.0, 1.1, and 1.2 and when  $DP_u$  is 0% MPa and  $T_u$  in the figure in the temperature range of 23.5–296.6 cm/s, 223.6–314.0 cm/s, 266.8–358.8 cm/s, and 301.0–399.6 cm/s, respectively, showing a trend increasing first and then decreasing. The maximum value occurs around  $\Phi$  of 1.1. In addition, from the high-temperature and low-pressure area to the low-temperature and high-pressure area,  $u_L$  shows a clear downward trend, and as the temperature increases, the pressure decreases and  $u_L$  changes faster, that is, a larger rate of change in  $u_L$  in the high-temperature and low-pressure area.

The feature points in the  $u_L$  range of 32.53–46.06 cm/s were further extracted to analyse the variation pattern of the  $u_L$  corresponding to  $\Phi$ . As illustrated in Figure 5, under a certain variation value of  $u_L$ , the corresponding  $P_u$  tends to increase first and then decrease, and the corresponding  $T_u$  decreases first and then increases as  $\Phi$  increases. The trend shows that, with the increase in  $\Phi$ , the isoline of  $u_L$  moves to the high-pressure and low-temperature region first and then moves to the low-pressure and high-temperature region around  $\Phi$  of 1.1. Figure 6 shows the rate of change  $R$  in  $u_L$  in the velocity range of 32.53–46.06 cm/s. It can be seen from the figure that  $R_{46.6-41.55} > R_{41.55-37.04} > R_{37.04-32.53}$ , which indicates that, in the range of  $\Phi$  at 0.9–1.2, the greater the  $u_L$ , the greater the rate of change  $R$  for  $u_L$ .

and 21.9–59.6 cm/s. With the increase in  $\Phi$ , the variations in  $u_L$  are 36.1 cm/s, 38.8 cm/s, 39 cm/s and 37.7 cm/s, respectively, showing a trend increasing first and then decreasing. The maximum value occurs around  $\Phi$  of 1.1. In addition, from the high-temperature and low-pressure area to the low-temperature and high-pressure area,  $u_L$  shows a clear downward trend, and as the temperature increases, the pressure decreases and  $u_L$  changes faster, that is, a larger rate of change in  $u_L$  in the high-temperature and low-pressure area.

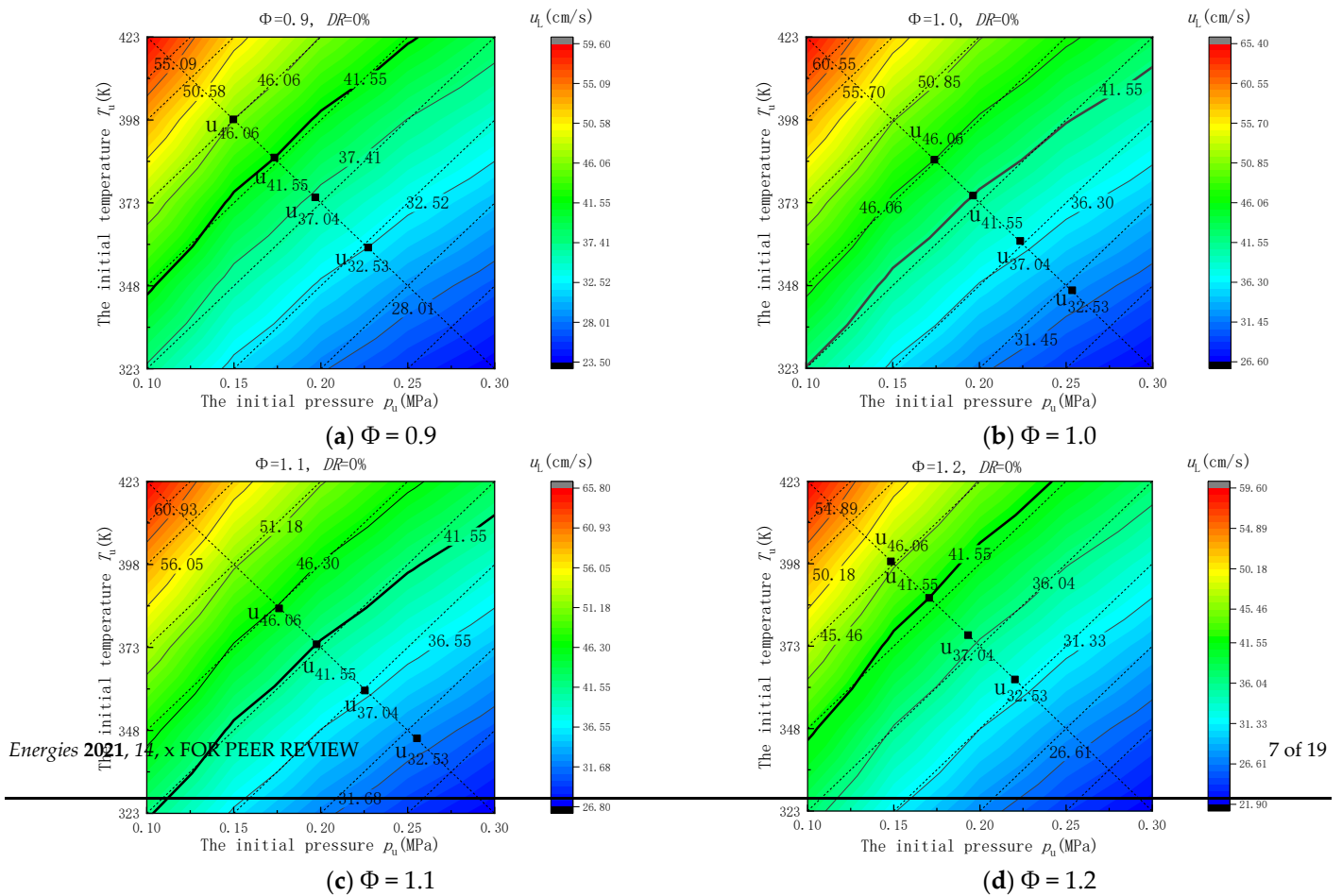


Figure 4. Equivalent distribution pattern of  $u_L$  under the coupling effect of  $T_u$  and  $p_u$  when the equivalence ratios are 0.9, 1.0, 1.1 and 1.2 and when dilution rate (DR) is 0%.

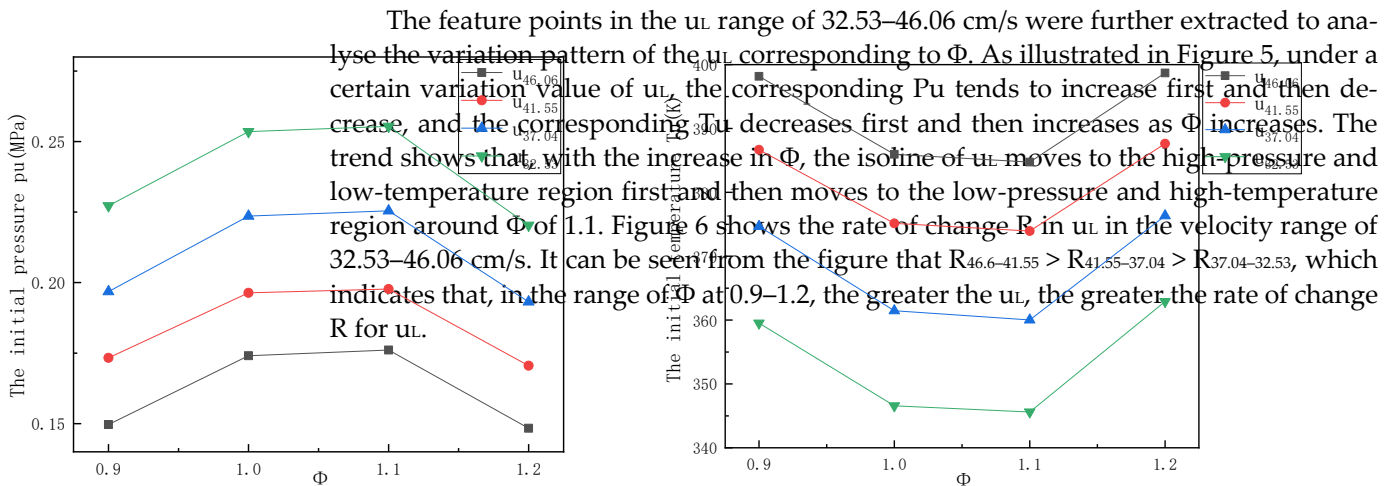


Figure 5. Equivalent variation of  $u_L$  corresponding to the initial condition change with  $\Phi$ .

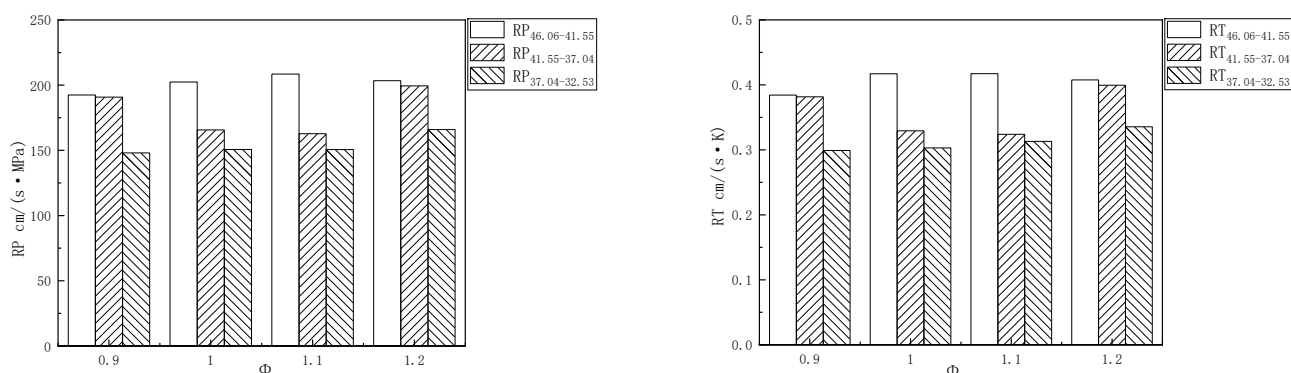


Figure 6. The rate of change  $R$  of  $u_L$  corresponding to  $\Phi$ .

The auxiliary dashed curves in Figure 4 stand for the equivalent effect on the  $u_L$  of  $T_u$  and  $P_u$ . The specific equivalent relationship is that the  $T_u$  change of  $\Delta T = 25 \text{ K}$  and the  $P_u$  change of  $\Delta P = 0.05 \text{ MPa}$  have equivalent effects on the same variation value of  $u_L$ . The intersection of the isoline of  $u_L$  and the auxiliary dashed curves in the figure indicates that  $T_u$  and  $P_u$  have an equal influence weight at this point. At this equal weight intersection point in the figure, the fracture area is fitted and fitted (reference Figure 7); Figure 8 shows the equivalent weight of  $u_L$ . Among them, the region II is the area where  $P_u$  has a greater influence and region I and region III are the areas where  $T_u$  has a greater influence. As  $\Phi$  increases, the equal weight influence of  $T_u$  and  $P_u$  exhibits different patterns. Referring to the figure, as  $\Phi$  increases, the  $u_L$  of the low-pressure equal weight curve ( $P_u = 0.15 \text{ MPa}$ ) gradually moves towards a higher  $P_u$ , but the increase rate changes more slowly with the increase in  $T_u$ . When  $T_u$  is greater than  $384 \text{ K}$  and  $\Phi$  is the low of the low-pressure equal weight curve reaches the maximum. As  $\Phi$  increases, the area of region II below  $384 \text{ K}$  gradually increases, that is, the area in which the initial temperature influence weight gradually increases. In the initial temperature above  $384 \text{ K}$ , as region I increases, the area of the region with the greater initial temperature influence weight increases first and then decreases. Furthermore, as  $\Phi$  increases and the initial temperature is greater than  $412 \text{ K}$ , the high-pressure equal weight curve ( $P_u = 0.22 \text{ MPa}$ ) gradually moves towards a higher  $P_u$ . When the initial temperature is in the range of  $342\text{--}412 \text{ K}$  and  $\Phi$  is 1.1, the  $P_u$  of the high-pressure equal weight curve reaches the minimum. When the initial temperature is lower than  $342 \text{ K}$  and  $\Phi$  is 1.2, the  $P_u$  of the high-pressure equal weight curve reaches the minimum. Therefore, as  $\Phi$  increases, the area in which the initial temperature is larger than  $412 \text{ K}$  with a greater initial temperature influence weight gradually decreases. While the initial temperature is below  $384 \text{ K}$ , the area with the greater initial temperature influence weight increases.

Figure 8 shows the equivalent distribution pattern of  $u_L$  under the coupling effect of initial temperature and  $P_u$  when  $\Phi$  is 1.0 and when DR is 0–16%. Under the temperature range of 323–423 K; the pressure range of 0.1–0.3 MPa; and the dilution rates of 0, 4%, 8%, 12%, 14% and 16%, the ranges of  $u_L$  are 26.6–65.4 cm/s, 19.1–49.4 cm/s, 13.5–37.2 cm/s, 9.35–27.7 cm/s, 7.7–23.75 cm/s and 6.25–20.3 cm/s. As DR increases, the variations in  $u_L$  are 38.8 cm/s, 30.3 cm/s, 23.7 cm/s, 18.35 cm/s, 16.05 cm/s and 14.05 cm/s, showing a gradually decreasing trend. It can be seen from the figure that  $u_L$  in the high-temperature and low-pressure region is still greater than that in the low-temperature and high-pressure region and that the increase in DR only affects the rate of change in  $u_L$ . That is, although the higher the  $u_L$ , the greater the rate of change in  $u_L$ , but the increase in DR will weaken the increase in the rate of change in  $u_L$ .



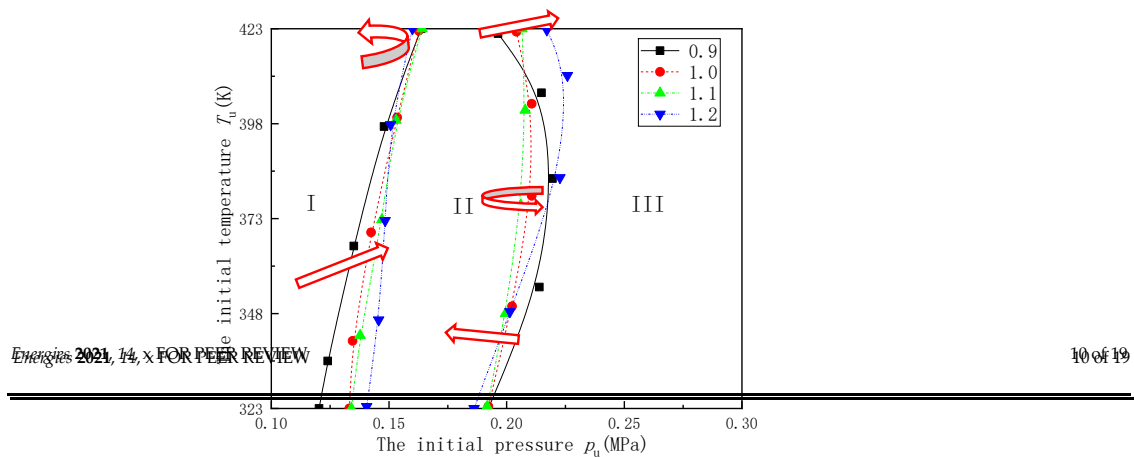


Figure 7. Variation pattern of the equal weight curve corresponding to  $\Phi$  with  $\Delta T \equiv 25$  K and  $\Delta P = 0.05$  MPa.

Figure 8 shows the equivalent distribution pattern of  $u_1$  under the coupling effect of initial temperature and  $P_u$  when  $\Phi$  is 1.0 and when DR is 0–16%. Under the temperature range of 323–423 K; the pressure range of 0.10–0.30 MPa, the variation rates of  $u_1$  are 0%, 4%, 8%, 12%, 14% and 16%. The ranges of  $u_1$  are 30.55–60.55 cm/s, 22.89–41.83 cm/s, 16.46–31.28 cm/s, 11.71–25.41 cm/s, 9.706–24.74 cm/s, and 8.006–18.54 cm/s, respectively. It can be seen from the figure that the variations in  $u_1$  are gradually increasing with increasing temperature and low-pressure region, and the high-pressure region is still greater than that in the low-pressure region. Although the increase in DR only affects the rate of change of  $u_1$ , it will weaken the rate of change in  $u_1$ .

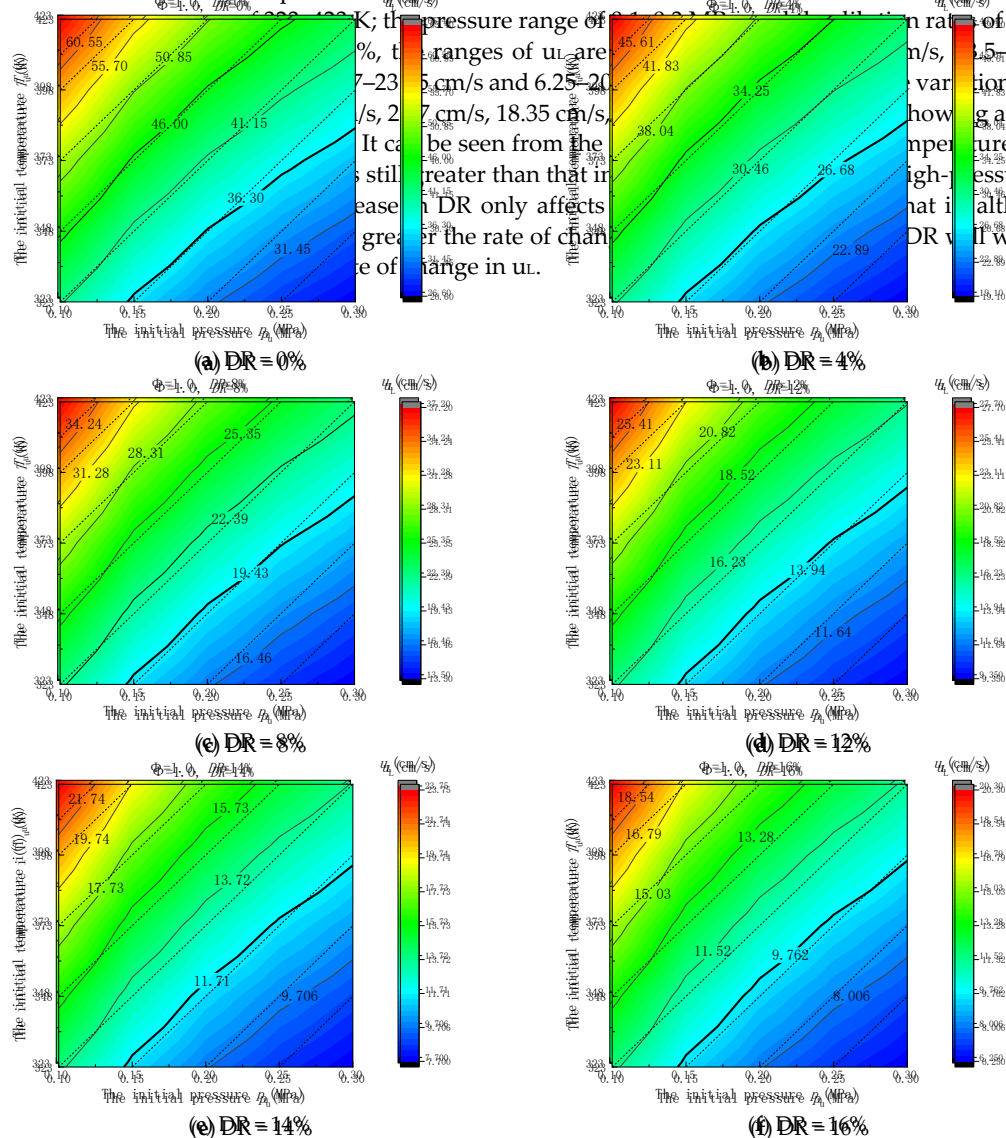
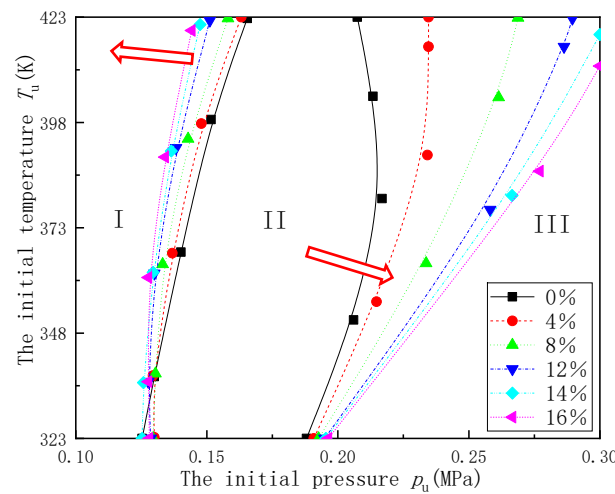


Figure 8. The equivalent distribution pattern of  $u_1$  under the coupling effect of  $T_u$  and  $P_u$  when  $\Phi$  is 1.0 and when the dilution rates are 0%, 4%, 8%, 12%, 14% and 16%.

Figure 9 shows the equal weight curve changes with the DR under the coupling effect of initial temperature and Pu. As illustrated with the increase in DR, the coupling effect of initial temperature and Pu is gradually weakened. When the initial temperature is greater than 350 K, the low-pressure equal weight curve (Pu = 0.125–0.165 MPa) gradually moves to the low-pressure area, and the equal weight curve (Pu = 0.125–0.165 MPa) gradually moves to the high-pressure area. When the initial temperature is 350 K and DR is 14%, Pu of the low-pressure equal weight curve increases as the area of region II also increases. When the initial temperature is 350 K and DR increases by 5%, the area of region II also increases, the initial temperature of 350 K gradually increases, and the equal weight curve also decreases. Additionally, with the increase in DR, the high-pressure equal weight curve (Pu = 0.18–0.25 MPa) gradually moves to the high-pressure area and the increase rate changes more steeply with the increase in initial temperature. As such, with the increase in DR, the area with the greater initial temperature influence and the greater rate of increase and the decrease rate gradually decreases.



**Figure 9.** Variation pattern of the equal weight curve corresponding to DR with  $\Delta T = 25$  K and  $\Delta P = 0.05$  MPa.

Figure 10 shows the equivalent distribution pattern of  $u_L$  under the coupling effect of  $T_u$  and DR when  $\Phi$  is 0.10 and when  $P_u$  is 0.25 MPa. The figure indicates that  $u_L$  gradually decreases from the high-temperature and low-dilution area to the low-temperature and high-dilution area. When  $P_u$  were 0.1 MPa, 0.15 MPa, 0.2 MPa, and 0.25 MPa, the range of  $u_L$  were 22.6–65.4 and 21.8–56.9 cm/s, 19.0–56.9 cm/s, and 14.9–54.6 cm/s. With the increase in  $P_u$ , the variation range of  $u_L$  were 4.2 cm/s, 3.9 cm/s, 3.5 cm/s, and 3.8 cm/s, respectively, decreasing trend. It can be seen that the larger the  $P_u$  is, the smaller the variation range of  $u_L$  is. The figure shows that in the parallel flow, the flow rate is relatively stable, and the relationship and hardly changes with the variation in  $P_u$ . The measurement shows that as the DR increases by 2% and decreases by 1% and the initial temperature (or decrease) by 29 decreases by 29 K, can be obtained.

The feature points in the  $u_L$  range of 29.54–40.82 cm/s were further extracted to analyse the variation pattern of  $u_L$  corresponding to  $P_u$ . As shown in Figure 11, when the variation in  $u_L$  is constant, the corresponding DR decreases with the increase in  $P_u$  and the initial temperature increases with the increase in  $P_u$ . This clearly shows that, with the increase in  $P_u$ , the isoline of  $u_L$  moves to the low-dilution rate and high-temperature area. Figure 12 shows the rate of change  $R$  of  $u_L$  in the velocity range of 29.54–40.82 cm/s. It can be seen from the figure that  $R_{40.82-37.06} > R_{37.06-33.30} > R_{33.30-29.54}$ , which indicates that, in the range of  $\Phi$  0.1–0.25, the greater the  $u_L$ , the greater the rate of change  $R$  in  $u_L$ .

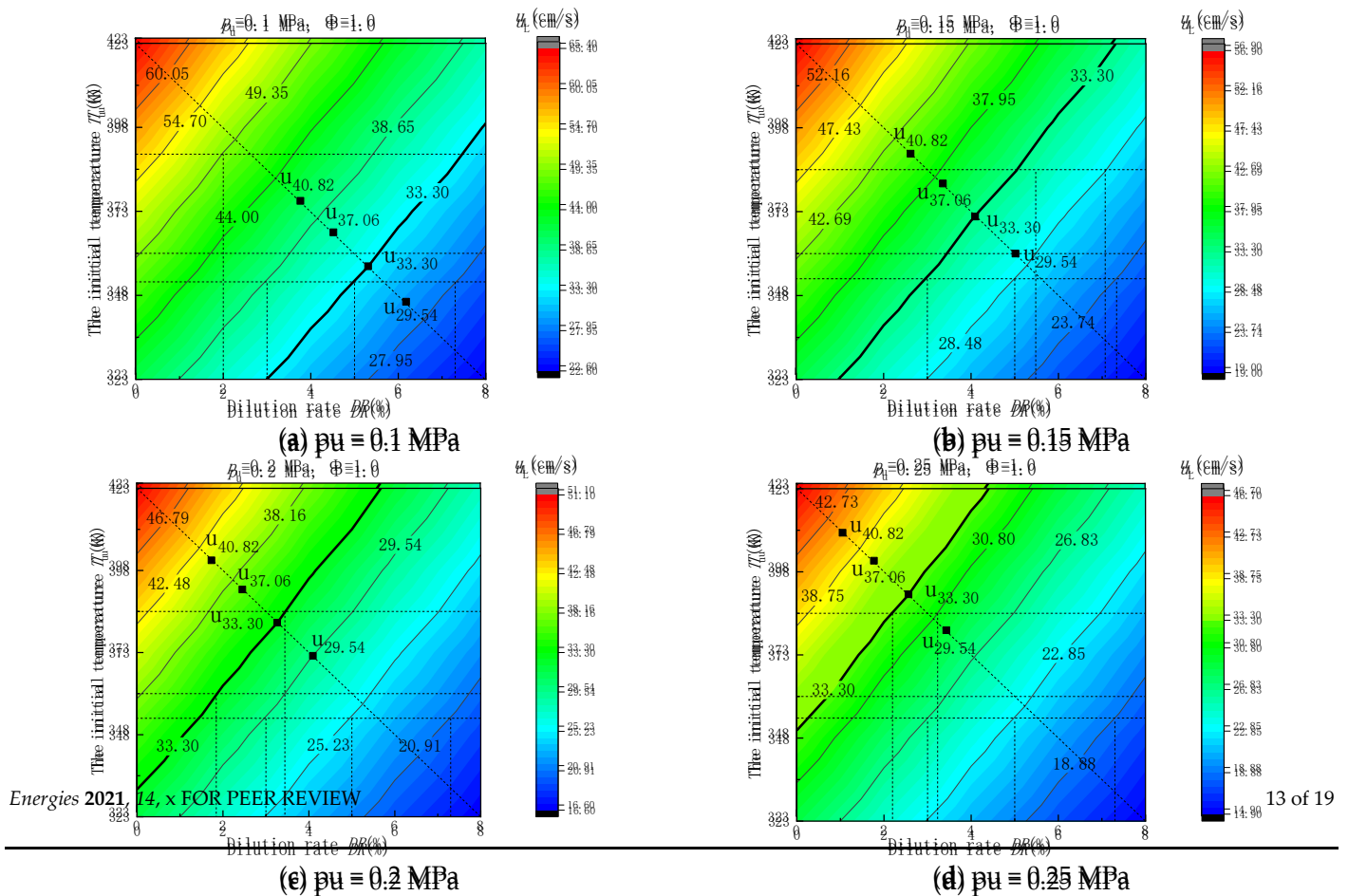


Figure 10. Equivalent distribution pattern of  $u_L$  under the coupling effect of  $T_{i0}$  and DR when  $\Phi$  is 1.0 and when the initial pressures are 0.1 MPa, 0.15 MPa, 0.2 MPa and 0.25 MPa.

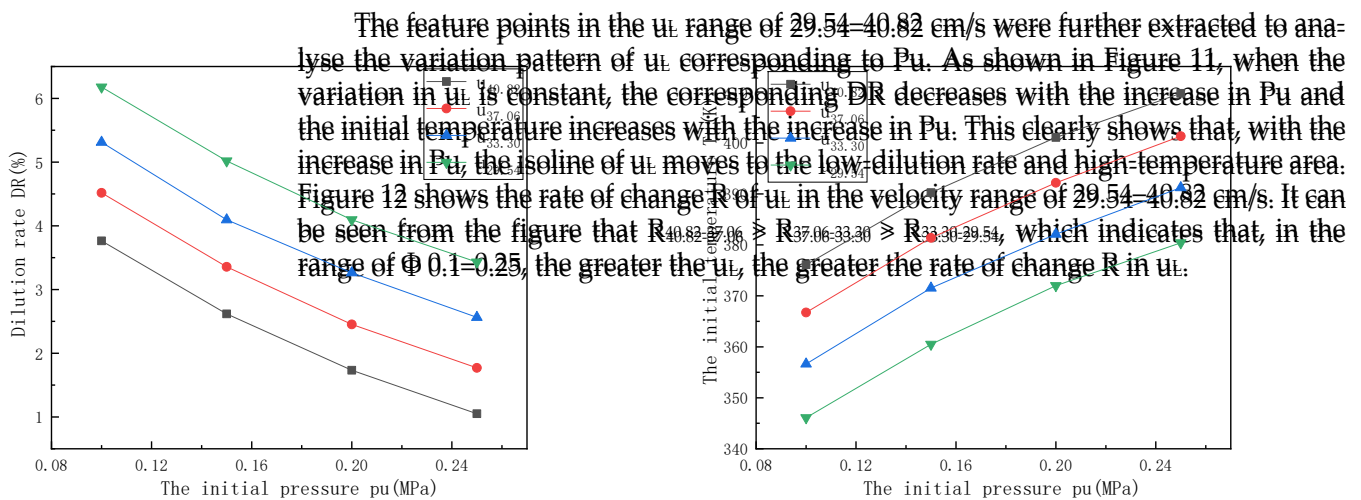
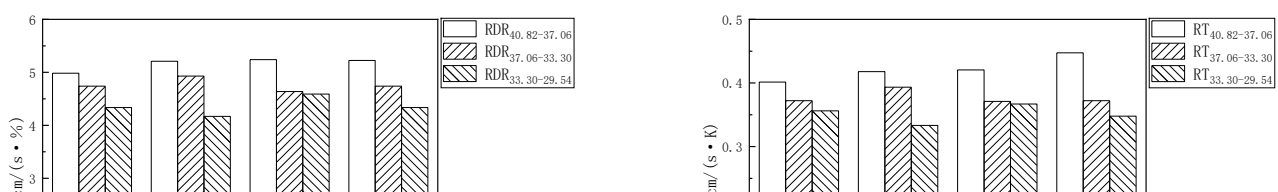


Figure 11. Equivalent variation of  $u_L$  corresponding to the initial condition change with  $p_u$ .



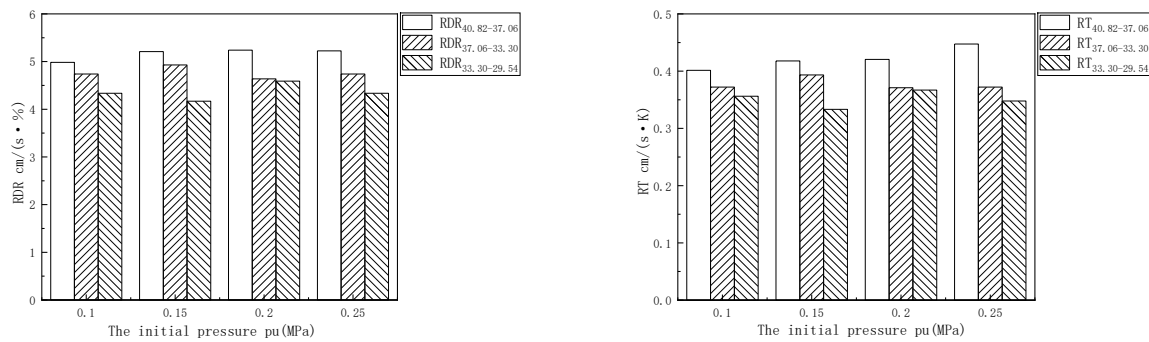
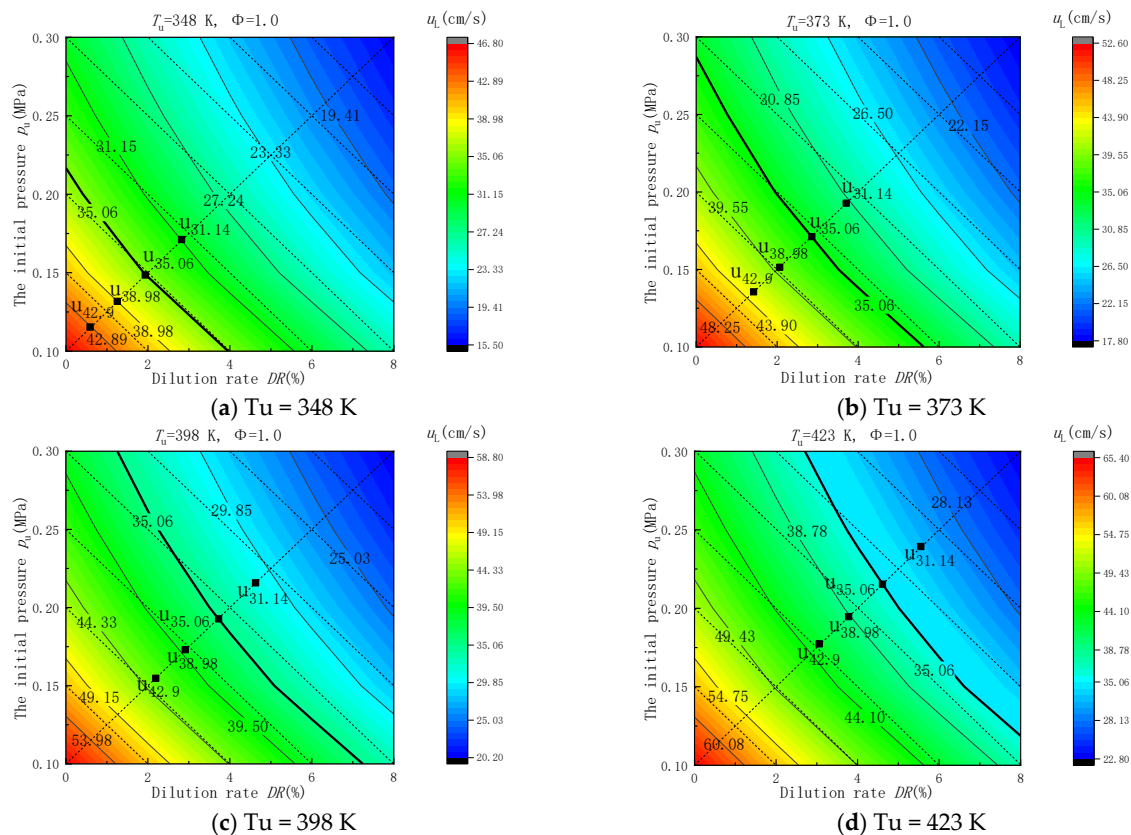
**Figure 11.** Equivalent variation of  $u_L$  corresponding to the initial condition change with  $P_u$ .**Figure 12.** The rate of change R of the  $u_L$  corresponding to the  $P_u$ .

Figure 13 shows the equivalent distribution pattern of  $u_L$  under the coupling effect of  $T_u$  and DR when  $\Phi$  is 1.0 and when the range of the initial temperature is 348–423 K. The figure indicates that  $u_L$  gradually decreases from the low-pressure and low-dilution rate area to the high-pressure and high-dilution rate area. When the initial temperatures are 348 K, 373 K, 398 K and 423 K, the ranges of  $u_L$  are 15.50–46.80 cm/s, 17.80–52.60 cm/s, 20.2–58.8 cm/s and 22.8–65.4 cm/s. With the increase in initial temperature, the variations in  $u_L$  are increasing trend. The larger the initial temperature,  $u_L$  respectively exhibiting a gradually increasing trend. The larger the initial temperature, the higher pressure and high dilution rate area, as well as the rate of change R in  $u_L$ . In addition, from the high-pressure and high-dilution rate area to the low-pressure and low-dilution rate area, a larger rate of change in  $u_L$  as  $P_u$  and DR increase,  $u_L$  changes faster, that is, a larger rate of change in  $u_L$  in the low-pressure and low-dilution rate area.

**Figure 13.** Equivalent distribution pattern of  $u_L$  under the coupling effect of  $P_u$  and DR when  $\Phi$  is 1.0 and when the initial temperatures are 348 K, 373 K, 398 K and 423 K.

The feature points in the  $u_L$  range of 31.14–42.9 cm/s were further extracted to analyse the variation pattern of  $u_L$  corresponding to  $P_u$ . As shown in Figure 14, when the variation in  $u_L$  is constant, the corresponding DR and  $P_u$  decrease with the increase in the initial temperature. This clearly shows that, with the increase in initial temperature, the isoline of  $u_L$  moves to the high-pressure and high-dilution rate area. Figure 15 shows the rate of change R in  $u_L$  in the velocity range of 31.14–42.9 cm/s. It can be seen from the figure that  $R_{398\text{K}} > R_{373\text{K}} > R_{348\text{K}}$ , which indicates that in the range of initial temperatures

The feature points in the  $u_L$  range of 31.14–42.9 cm/s were further extracted to analyse the variation pattern of  $u_L$  corresponding to  $P_u$ . As shown in Figure 14, when the variation in  $u_L$  is constant, the corresponding DR and  $P_u$  decrease with the increase in the initial temperature. This clearly shows that, with the increase in initial temperature, the isoline of  $u_L$  moves to the high-pressure and high-dilution rate area. Figure 15 shows the rate

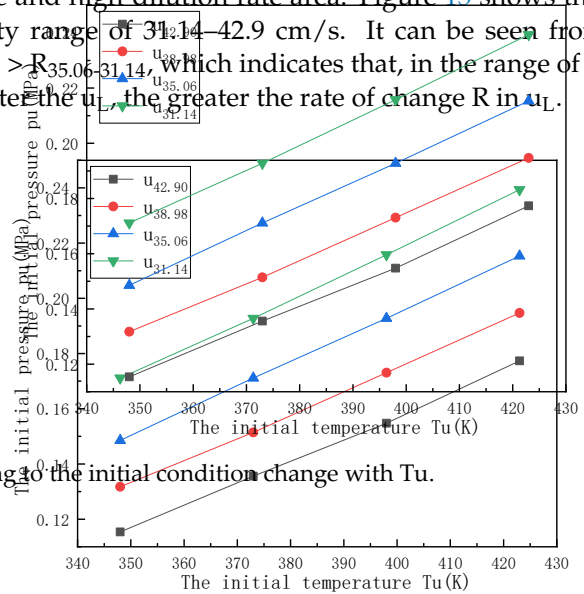
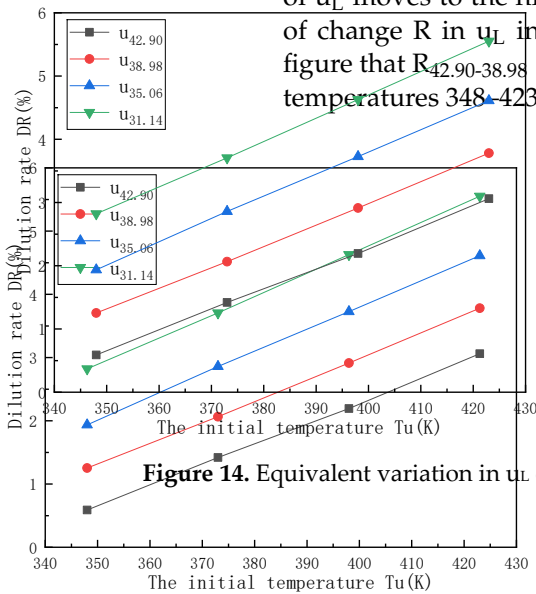


Figure 14. Equivalent variation in  $u_L$  corresponding to the initial condition change with  $T_u$ .

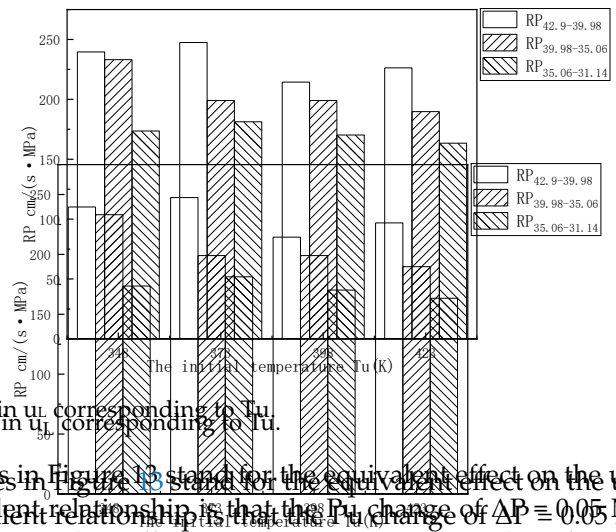
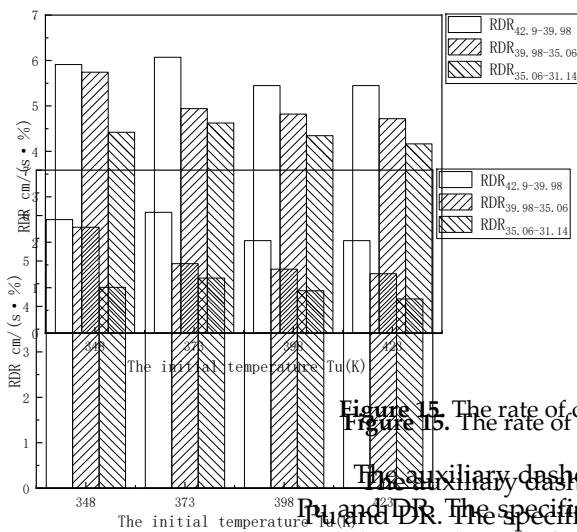
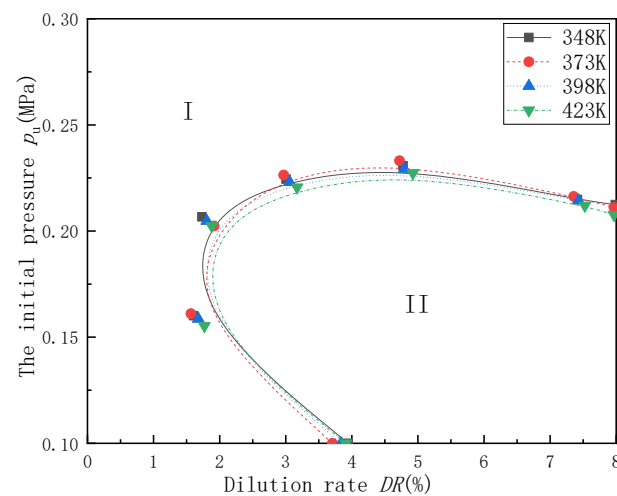


Figure 15. The rate of change R in  $u_L$  corresponding to  $T_u$ .

The auxiliary dashed curves in Figure 13 stand for the equivalent effect on the  $u_L$  of  $P_u$  and DR. The specific equivalent relationship is that the  $P_u$  change of  $\Delta P = 0.05$  MPa and the DR change of 2% have equivalent effects on the same variation value of  $u_L$ . The intersection of the isolines of  $u_L$  and the auxiliary dashed curves in the figure indicates that  $P_u$  and DR have an equal influence weight on  $u_L$  at this point. All the equal weight intersection points in the figure are extracted and fitted onto curves (refer to Figure 16), and two regions are formed by equal weight curves. Region I is the area where the DR has a greater influence on  $u_L$ , while in region II,  $P_u$  has a greater influence on  $u_L$ . This shows that, under the influence of the coupling effect of  $P_u$  and DR on  $u_L$ ,  $P_u$  of the low-pressure and high-dilution rate area has a greater weight on  $u_L$ . As the initial temperature increases, the equal weight curves of  $P_u$  and DR exhibit different movement patterns. Region I is the area where the DR has a greater influence on  $u_L$ , while in region II,  $P_u$  has a greater influence on  $u_L$ . This shows that, under the influence of the coupling effect of  $P_u$  and DR on  $u_L$ ,  $P_u$  of the low-pressure and high-dilution rate area has a greater weight on  $u_L$ . As the initial temperature increases, the equal weight curves of  $P_u$  and DR exhibit different movement patterns.



**Figure 16.** Variation pattern of the equal weight curve corresponding to  $T_u$  with  $\Delta DR 2\%/2\%$  and  $P_i = 0.05$  MPa.

## 5. Conclusions

Due to the initial temperature having a promoting effect on laminar burning velocity while the initial pressure and dilution rate have an inhibitory effect on laminar burning velocity, equal laminar burning velocities can be obtained by initial condition coupling with different temperatures, pressures and dilution rates. This paper analysed the equivalent distribution pattern of the laminar burning velocity and the variation pattern of the equal weight curve by coupling effect of the initial pressure (0.1–0.3 MPa), initial temperature (323–423 K) and dilution rate (0–16%). The main conclusions are summarized as follows:

- As the initial temperature increases and the initial pressure decreases, the rate of change in laminar burning velocity increases. Moreover, the increase in equivalent ratio makes the variation of laminar burning velocity show a trend increasing first and then decreasing and the maximum value is reached when  $\Phi = 1.1$ , while with the increase in the dilution rate, the variation in laminar burning velocity gradually decreases and the maximum value is reached when  $\Phi = 1.1$ , while with the increase in the dilution rate, the variation in laminar burning velocity gradually decreases and the isoline of laminar burning velocity gradually moves towards the high-temperature and low-pressure area.
- The equivalent effect of initial condition coupling can obtain equal laminar burning velocity with the dilution rate increasing (or decreasing) by 2% and the initial temperature increasing (or decreasing) by 29 K. Additionally, the variation in laminar burning velocity tends to decrease with the increase in initial pressure, and the isoline of laminar burning velocity gradually moves towards the high-temperature and low-dilution rate area.
- With the dilution rate and the initial pressure decreasing, the rate of change in laminar burning velocity increases. Additionally, the variation in laminar burning velocity tends to increase with the increase in initial temperature, and the isoline of laminar burning velocity gradually moves towards the high-temperature and low-dilution rate area.
- As the equivalent ratio increases additionally, the temperature is less than 384 K, the area with the greater initial temperature influence weight gradually increases, and when it exceeds 384 K, the area gradually increases first and then decreases, and the high-temperature area gradually decreases with the increase in equivalence ratio, while when the initial temperature is less than 342 K, the area gradually increases with the increase in equivalence ratio.
- With the increase in dilution rate, the area with the greater initial temperature influence weight gradually decreases and the decrease rate gradually decreases. In the low-pressure and high-dilution rate region, the initial pressure has a greater influence weight on laminar burning velocity, while in the high-pressure and low-dilution rate region, the dilution rate has a greater influence weight on laminar burning velocity.

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## Abbreviations

Pu	Initial pressure, MPa
Tu	Initial temperature, K
$u_L$	Laminar burning velocity, $m \cdot s^{-1}$
DR	Dilution rate
$\Phi$	Equivalence ratio
CVC	Constant volume chamber
EGR	Exhaust gas recirculation
$R_u$	Instantaneous flame radius: mm
$S_n$	Stretched flame propagation speed, $m \cdot s^{-1}$
$\kappa$	Curvature of spherical flame, $mm^{-1}$
K	Stretch rate, $s^{-1}$
$L_b$	Markstein length, mm
$S_l$	Unstretched flame propagation speed, $m \cdot s^{-1}$
Le	Lewis number
$\lambda$	Thermal conductivity, $W \cdot (m \cdot K)^{-1}$
$\rho_u$	Density of unburned gas, $kg \cdot m^{-3}$
$c_p$	Specific heat capacity, $J \cdot (kg \cdot K)^{-1}$
$D_m$	Mass diffusion coefficient, $m^2 \cdot s^{-1}$
$D_T$	Thermal diffusion coefficient, $m^2 \cdot s^{-1}$
$\rho_b$	Density of burned gas, $kg \cdot m^{-3}$
$\sigma$	Thermal expansion ratio

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