

## Energy harvesting using a magnetostrictive transducer based on switching control

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### ABSTRACT

In this work, a switching control energy harvesting method using magnetostrictive materials is proposed. By combining a magnetostrictive material, an electric circuit, and an electronic switch, large-scale kinetic to electrical energy conversion can be achieved. The magnetostrictive material, magnet bias, and coils constitute an energy transducer, called a magnetostrictive transducer. The electronic switch strategically controls the switching of the circuit state according to an input switching signal. Using numerical simulations, we optimised the parameters and validated the harvesting performance with experimental measurements using a 3.75 m vibrated cantilever truss structure. In 20.0 s, the proposed method achieved an electrical energy of approximately 45  $\mu$ J, which is seven times more than that of the conventional passive method.

## 1. Introduction

### 1.1. Background

Recently, wireless sensor networks (WSNs) [1] and Internet of Things (IoTs) [2] have attracted considerable attention and are being rapidly adopted in various fields. Energy harvesting is a promising technique that can produce renewable and clean electrical energy. Roadways and bridges are major infrastructures involved in daily life, and the movement of people and vehicles constantly affects these infrastructures through structural vibrations. The energy harvesting technique used for such systems [3] is schematically shown in Fig. 1. The vibration generated from moving transportation systems can provide a renewable and sustainable energy source that can potentially power structural health monitoring sensors.

Kinetic to electrical energy conversion can be achieved using various smart materials, including piezoelectric materials [4,5], dielectric elastomers [6], shape memory alloys [7], electrorheological fluids [8], magnetorheological fluids [9], magneto rheological elastomers [10,11], and magnetostrictive materials [12,13].

The energy harvesting method used in piezoelectric materials has

been widely studied [14]. Lombardi and Lallart [15] proposed a synchronous electric charge and induced current extraction (SECICE) technique that can significantly increase output power, particularly for low-coupled or highly damped systems. For switching control to amplify energy harvesting, synchronised switch harvesting on an inductor (SSHI) has been proposed to control switch devices according to vibration displacements [16]. Jia and Seshia [17] presented a power optimisation design for piezoelectric cantilever vibration energy harvesters based on experimentally validated analytical and numerical analysis. Hara et al. [18–20] developed a self-sensing estimation method for switching control energy harvesters based on the SSHI method and proposed novel switching strategies for enhancing energy harvesting by addressing the vibration suppression effect. However, piezoelectric materials mainly comprise brittle ceramics, which may be unsuitable for complex or extreme environments. Hence, magnetostrictive materials have gained attention in recent years; they are a type of smart material comprised primarily of iron and rare-Earth elements, such as Terfenol-D (Tb-Dy-Fe), Fe-Co alloys, and Galfenol (Fe<sub>81.4</sub>Ga<sub>18.6</sub>) [21]. Because these materials are iron alloys, they have high mechanical strength and workability, which make them suitable for various complex environments.

Magnetostrictive materials can be used for conversion between

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### Nomenclature

$b_m$ :	magnetostriction coefficient of magnetostrictive materials [V·s/m = N/A]
$\mathbf{B}_m$ :	magnetostriction coefficient matrix
$C_i$ :	capacitance of inversion capacitor [F]
$C_h$ :	capacitance of harvesting capacitor [F]
CPCS:	current peak control strategy
$\mathbf{D}$ :	damping coefficient of structure [N·s/m]
$E_h$ :	harvested electrical energy [J]
$\mathbf{f}_{\text{ext}}$ :	multiple external force vector
$F_m$ :	external force of magnetostrictive transducer [N]
$I$ :	current [A]
$\mathbf{I}$ :	induced current vector of magnetostrictive transducers
$k_m^I$ :	stiffness of magnetostrictive transducer at constant current [N/m]
$k_{ii}$ :	stiffness of $i$ -th bar member.
$\mathbf{K}$ :	constant-current stiffness matrix of structure
$L_m^S$ :	inductance of magnetostrictive transducer at constant

	velocity [H]
$\mathbf{M}$ :	mass matrix of structure
$P_i$ :	axial force applied to $i$ -th bar member [N]
$Q_i$ :	charge of inversion capacitor [C]
$Q_h$ :	charge of harvesting capacitor [C]
$R_0$ :	internal resistance of coil [ $\Omega$ ]
$T$ :	kinetic energy [J]
$u$ :	elongation of magnetostrictive materials
$V_m$ :	induced voltage of magnetostrictive transducer [V]
$V_i$ :	voltage of $C_i$ [V]
$V_h$ :	harvested voltage [V]
VPCS:	velocity peak control strategy
$W$ :	potential energy [J]
$x$ :	displacement of structure [m]
$\mathbf{x}$ :	position coordinates of each node
$\zeta_e$ :	damping coefficient of inductor-capacitor electrical oscillation
$\omega_e$ :	frequency of inductor-capacitor electrical oscillation
$\Delta t$ :	time difference [s]

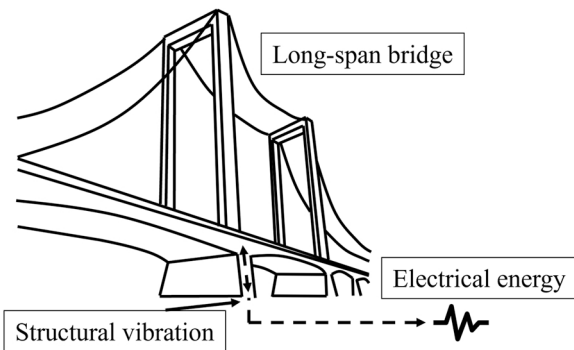


Fig. 1. Schematic of energy harvesting used in roadways and bridges [3].

kinetic and magnetic energy, which is referred to as the Villari-Joule effect [12]. The Joule effect is a property of ferromagnetic materials that causes them to generate a mechanical strain when subjected to a magnetic field. Meanwhile, the Villari effect is a property of magnetostrictive materials that causes them to become magnetically susceptible when subjected to mechanical stress. Magnetostrictive materials are typically used to design sensors and actuators [22–24]. Viola et al. [25] proposed a power generator using Terfenol-D for energy harvesting from traffic vibrations and experimentally validated this system. Zucca et al. [26] analysed the quantities that influence the performance of a direct force transducer based on Terfenol-D. The power generated by the transducer has a complex dependence on the coil characteristics, type of permanent magnets used, and mechanical excitation characteristics. Mori et al. [27] constructed a giant magnetostrictive cantilever with a Terfenol-D layer, stainless steel layer, and movable proof mass designed to automatically adjust the resonant frequency to match the external vibration frequency in real-time. Researchers focused on the use of Galfenol and found that the magnetostrictive effect of Galfenol was optimal when the composition ratio of Ga was approximately 20% [28–30]. They also clarified the influence of an external magnetic field on the magnetostrictive effect by focusing on the elastic modulus of the magnetostrictive material under the influence of an optimal external magnetic field. Kita et al. [31] developed a high-power magnetostrictive vibration power generator for battery-free wireless electronics based on a cantilever parallel beam structure consisting of coil-wound Galfenol and stainless plates with a permanent magnet for bias. Ueno [32]

proposed a simple, robust, and highly sensitive vibration power generation device based on the unimorph of an Fe-Ga plate and U-shaped frame. Li et al. [33] presented an energy harvester employing a cantilever beam and magnetostrictive/piezoelectric laminate transducer to transform rotational energy into electrical energy. Li et al. [34] proposed a ferro-nickel (Fe-Ni)/PZT H-type fork magneto-electric composite structure with a high quality value and energy management circuit for harvesting. Iannone et al. [35] connected the magnetostrictive harvester to an AC-DC boost converter. Based on the simulation results, the proposed circuit was proven to increase more output voltage than that induced by the magnetostrictive harvester. Clemente et al. [36] presented the experimental verification of an AC-DC boost converter driven with a real-time operating Arduino board.

### 1.2. Objectives

Combining energy harvesting control and piezoelectric materials has attracted considerable research interest. Correspondingly, improvements in magnetostrictive material properties have been a research focus, while several studies have focused on energy harvesting control to achieve more efficient harvesting performance. Conventional studies typically utilise a half rectifier circuit to passively harvest electrical energy.

The objective of this study is to propose a novel energy harvesting method using a magnetostrictive transducer. The method differs from that of conventional studies because strategical switching between multiple electrical systems is utilised. This results in a better energy harvesting performance because the properties of the raw materials used are not altered, which is achieved by combining resistor-inductor (RL) and resistor-inductor-capacitor (RLC) circuits with a magnetostrictive transducer. As the state of the circuit changes, the electrical energy temporarily stored in the coils and capacitor is amplified, thereby amplifying the harvested electrical energy. Subsequently, two control strategies that can be synchronised using mechanical vibrations to switch between the circuit states are discussed.

The proposed switching control strategies are applied to a truss structure to simulate flexible infrastructure. The harvesting performance of the proposed method is predicted by numerical simulations and validated by experimental measurements.

## 2. Theoretical analysis, circuit design, and switching control strategies

### 2.1. Equation of magnetostrictive transducer

The overall strength when the transducer is embedded in the structure must be considered. Although the magnetostriction of Terfenol-D is the largest, it cannot play a load-bearing role owing to its fragility. While ensuring relatively large magnetostriction, the mechanical strength of Galfenol is higher than that of Terfenol-D. Considering the combination with target structures, we assumed that the magnetostrictive materials and other components exhibit similar Young's modulus. Hence, Galfenol ( $\text{Fe}_{81.4}\text{Ga}_{18.6}$ ) was selected as the magnetostrictive material.

The magnetostrictive transducer used consists of Galfenol, coils, and neodymium magnets, as shown in Fig. 2(a). Magnetostrictive transducers are an electromechanically coupled system, and their fundamental equations can be derived as

$$F_m = k_m^I u - b_m I, \quad (1)$$

$$V_m = -b_m \dot{u} - L_m^S \dot{I} - R_0 I, \quad (2)$$

where  $k_m^I$  is the stiffness of the magnetostrictive transducer at constant current,  $L_m^S$  is the inductance of the magnetostrictive transducer at constant velocity,  $b_m$  represents the magnetostriction coefficient of the magnetostrictive transducer, while  $F_m$ ,  $V_m$ ,  $u$ , and  $I$  represent the external force acting on the magnetostrictive transducer, induced voltage of the magnetostrictive transducer, elongation of the magnetostrictive material, and current, respectively. Therefore, by using a magnetostrictive transducer, the kinetic energy from vibrations can be converted into electrical energy. Magnetostrictive transducers behave as a voltage source in tandem with the coils and an internal resistor  $R_0$ . Fig. 2(b) shows the schematic of an electrical model for such systems.

### 2.2. Proposed control circuit and mechanism

Fig. 3 illustrates a conventional passive harvesting system in which the magnetostrictive transducer is assembled in tandem with a diode to form a half rectifier passive circuit. Because the conventional passive harvesting system using a diode only aims to rectify the AC voltage  $V_m$  to a DC voltage, the maximum value of the DC voltage will not exceed the amplitude of  $V_m$ . Hence, we consider that this diode-based passive harvesting method will not amplify the harvested energy. However, this passive method cannot amplify harvested energy from the magnetostrictive transducer.

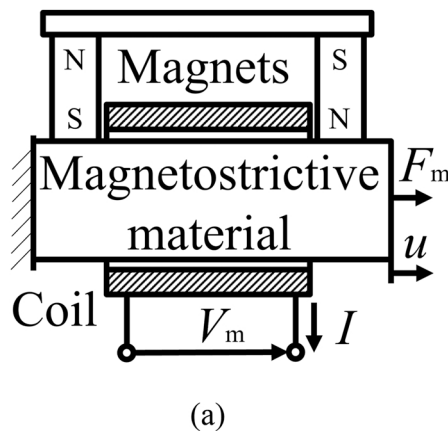


Fig. 2. Schematic of magnetostrictive transducer and corresponding electrical circuit:  $L_m^S$ , inductance of magnetostrictive transducer at constant velocity;  $b_m$ , magnetostriction coefficient of magnetostrictive transducer;  $F_m$ , external force acting on magnetostrictive transducer;  $V_m$ , induced voltage of magnetostrictive transducer;  $u$ , elongation of the magnetostrictive material;  $I$ , current;  $R_0$ , internal resistor.

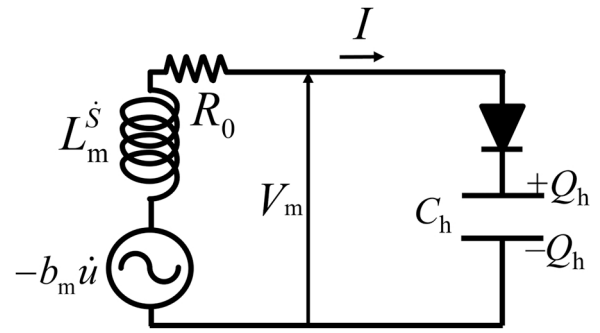


Fig. 3. Electrical model of magnetostrictive transducer assembled with half rectifier passive circuit:  $L_m^S$ , inductance of magnetostrictive transducer at constant velocity;  $b_m$ , magnetostriction coefficient of magnetostrictive transducer;  $V_m$ , induced voltage of magnetostrictive transducer;  $u$ , elongation of the magnetostrictive material;  $I$ , current;  $R_0$ , internal resistor;  $C_h$ , harvesting capacitor;  $Q_h$ , charge of harvesting capacitor.

The harvesting performance of magnetostrictive transducers is improved by introducing the circuit shown in Fig. 4, which consists of a magnetostrictive transducer, two capacitors ( $C_i$ ,  $C_h$ ), two diodes, and an electronic switch. Here,  $L_m^S$  is combined with the inversion capacitor  $C_i$  to form a high-frequency inductor-capacitor (LC) electrical oscillation.  $C_h$  is the harvesting capacitor that accumulates electrical energy. The electronic switch is selected between points X and Y. Fig. 5 shows each

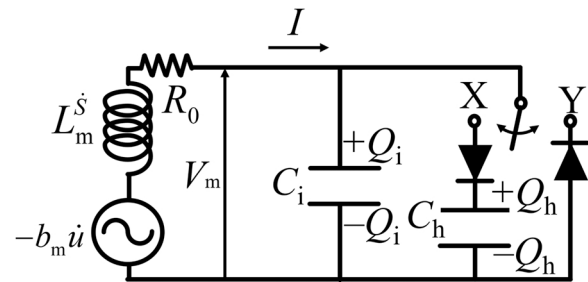


Fig. 4. Proposed circuit assembly for magnetostrictive transducer with electronic switch:  $L_m^S$ , inductance of magnetostrictive transducer at constant velocity;  $b_m$ , magnetostriction coefficient of magnetostrictive transducer;  $V_m$ , induced voltage of magnetostrictive transducer;  $u$ , elongation of the magnetostrictive material;  $I$ , current;  $R_0$ , internal resistor;  $C_h$ , harvesting capacitor;  $Q_h$ , charge of harvesting capacitor;  $Q_i$ , charge of inversion capacitor.

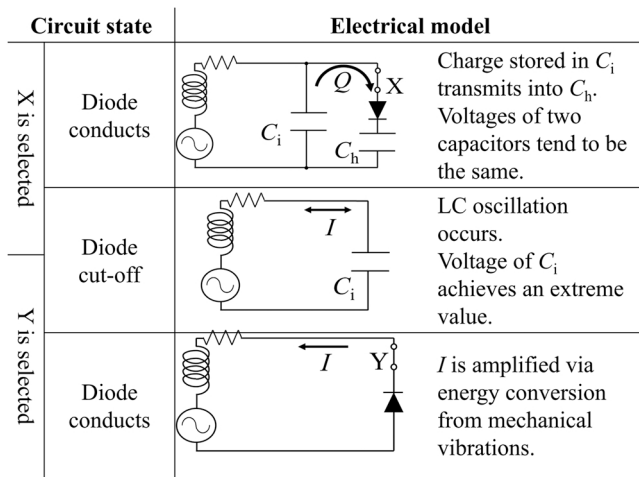


Fig. 5. Current and voltage mechanisms of proposed circuit:  $I$ , current;  $C_h$ , harvesting capacitor;  $C_i$ , inversion capacitor;  $Q$ , charge; LC, inductor-capacitor.

the magnetostrictive transducer will increase with the structural vibration.  $I$  can be calculated as

$$I = -\exp\left(-\frac{R_0}{L_m^s}t\right) \int \frac{b_m \dot{u}}{L_m^s} \exp\left(\frac{R_0}{L_m^s}t\right) dt. \quad (5)$$

where  $t$  is time. Because  $C_h$  and  $C_i$  are shorted, the harvesting process does not occur, and the electrical energy is temporarily stored in  $L_m^s$ .

Regardless of the selected point, if the diode is not conducting,  $I$  flows through  $C_i$ . The LC electrical oscillation caused by  $L_m^s$  and  $C_i$  occurs. The circuit equation can be described as

$$L_m^s \dot{I} + R_0 I + \frac{Q_i}{C_i} = -b_m \dot{u}, \quad \dot{Q}_h / C_h = 0. \quad (6)$$

Eq. (6) can be rewritten as

$$L_m^s \ddot{Q}_i + R_0 \dot{Q}_i + \frac{Q_i}{C_i} = -b_m U \omega \cos(\omega t + \gamma), \quad \dot{Q}_h / C_h = 0. \quad (7)$$

The solution of Eq. (7) can be determined as

$$Q_i = A \exp(-\zeta_e \omega_e t) \sin(\omega_e \sqrt{1 - \zeta_e^2} t + \theta - \lambda) + \frac{b_m U}{L_m \sqrt{(\omega_e^2 - \omega^2)^2 + 4\omega_e^2 \omega^2 \zeta_e^2}} \cos(\omega t + \gamma + \sigma), \quad (8)$$

circuit state during the switching and energy harvesting processes. The elongation  $u$  is assumed to be a sinusoidal function described by  $u = U \sin(\omega t + \gamma)$ . The proposed circuit contains three circuit states: short, LC, and charge transfer.

When point X is selected, if the diode is conducting, the harvesting process can be divided into passive and charge transfer processes. When the passive process occurs, the circuit connected to the anode of the diode can be considered to be an open circuit, and the input voltage of the diode  $V_D$  can be described as

$$V_D = -b_m \dot{u} - L_m^s \dot{I} - R_0 I. \quad (3)$$

$$V_i = \frac{A \exp(-\zeta_e \omega_e t)}{C_i} \sin(\omega_e \sqrt{1 - \zeta_e^2} t + \theta) + \frac{b_m U}{C_i L_m \sqrt{(\omega_e^2 - \omega^2)^2 + 4\omega_e^2 \omega^2 \zeta_e^2}} \cos(\omega t + \gamma + \sigma). \quad (10)$$

Because the magnetostrictive transducer is opened,  $I$  decreases to zero. When  $I = 0$ ,  $V_D$  approaches  $-b_m \dot{u}$ . The harvested voltage of  $C_h$  becomes infinitely close to  $V_D$  when the diodes are assumed to be ideal. Therefore, the passive method has a limit for energy harvesting.

When the charge transfer occurs, the charge stored in  $C_i$  transfers into  $C_h$  due to a potential difference. Increasing the potential difference between the two capacitors can increase the voltage of  $C_h$ . However, as in the passive process described above, the voltage of  $C_h$  cannot be larger than that of  $C_i$ .

When point Y is selected, the circuit connected in tandem forms a short circuit if the diode is conducting. The circuit equation can be described as

$$L_m^s \dot{I} + R_0 I = -b_m \dot{u}, \quad Q_i / C_i = 0, \quad \dot{Q}_h / C_h = 0, \quad (4)$$

where  $Q_i$  and  $Q_h$  are the charges of the inversion and harvesting capacitors, respectively. In this process, the electrical energy induced from

where  $A$  and  $\theta$  are constants determined by the initial conditions;  $\omega_e$  and  $\zeta_e$  represent the frequency and damping coefficient of LC electrical oscillation, respectively; and  $\omega_e$ ,  $\zeta_e$ ,  $\lambda$ , and  $\sigma$  can be defined as

$$\omega_e \equiv \frac{1}{\sqrt{L_m^s C_i}}, \quad \zeta_e \equiv \frac{R_0}{2} \sqrt{\frac{C_i}{L_m^s}}, \quad \lambda \equiv \tan^{-1}\left(\frac{\sqrt{1 - \zeta_e^2}}{\zeta_e}\right), \quad \sigma \equiv \tan^{-1}\left(\frac{2\omega_e \omega \zeta_e}{\omega_e^2 - \omega^2}\right) \quad (9)$$

The voltage of  $C_i$  can be summarised as

$V_i$  is determined by the initial conditions of the circuit state before switching. Therefore, the time associated with changing circuit states determines the maximum  $V_i$  value. Eq. (10) comprises electrical and mechanical terms. The first term refers to the voltage of LC electrical free oscillation, while the second term refers to the induced voltage from structural vibration. Because the amplitude of the first term is related to time  $t$ , the amplitude is considerably affected by the exponential function. In contrast, the amplitude of the second term is constant. Furthermore,  $C_i$  is selected as an extremely small value to ensure that the LC oscillation frequency is considerably higher than the structural vibration frequency. During the LC oscillation, the first term is more active with time  $t$ . Therefore, this term has a greater effect on  $V_i$ .

### 2.3. Theoretical control strategies

In conventional piezoelectric energy harvesting technologies, LC oscillation is used to amplify the induced voltage converted from

mechanical vibrations. These technologies have associated switching control strategies for synchronising electromechanical systems. These conventional strategies for the piezoelectric elements work by detecting the peaks of the mechanical displacement or electrical charge. In the case of magnetostrictive transducers, we focused on the induced current since the magnetostrictive transducer was assimilated into the inductor as shown in Fig. 2(b). Because the induced current is related to the strain velocity, as described in Eq. (2), the control strategies for the magnetostrictive transducer should also focus on the strain velocity or the induced current. Herein, two switching control strategies are developed.

Within the mechanical system, when the vibration velocity reaches the maximum value, the kinetic energy is high because the magnetostrictive transducer induces the most electrical energy when the kinetic energy of the mechanical system is maximised. Hence, the velocity peak control strategy (VPCS) is proposed. VPCS aims to switch between connection points when the structural vibration velocity is maximised. To realise VPCS, the velocity of the structural vibration is required.

Within the electrical system, when the induced current reaches a maximum value, the harvested energy through the electrical circuit is high because internal resistance is inevitable in an actual electrical circuit. The composition of the internal resistance and inductance of the magnetostrictive transducer cause a phase shift between the velocity and induced current. Hence, the current peak control strategy (CPCS) is also proposed. CPCS aims to switch between connection points when the induced current is maximised. To realise CPCS, the induced current of the electrical circuit is required.

#### 2.4. Applicable control strategies using truss displacement

Using a single-mode vibration application environment, we confirm that the acceleration of the elongation and displacement of the structure maintain the same phase. Considering the compositions of the magnetostrictive transducer and truss structure, we chose the displacement of the structure to predict the velocity peak. Therefore, we proposed VPCS as

[VPCS]:  
When  $x > 0$ , point Y should be selected.  
When  $x < 0$ , point X should be selected.

Here,  $x$  is the displacement of the structure.

Regarding CPCS, we need to detect the induced current of the electrical circuit. However, due to noise and LC electrical oscillation, the current peak is difficult to detect. Because the initial magnetic field is considerably larger than the magnetic field changing due to strain vibrations in our experimental environment, we assumed that this magnetostrictive transducer shows a linear relationship between the magnetic field and strain. Furthermore, under this initial magnetic field, the inductance and internal resistance maintain constant values. The phase difference between the vibration velocity and induced current is determined by the inductance and resistance. Hence, we designed CPCS based on VPCS and the phase difference between the vibration velocity and circuit current. When we only consider the case of the first natural vibration mode, the current peak can be calculated by introducing a phase difference based on the vibration velocity; we propose the following method to simulate CPCS. An appropriate time based on the vibration velocity and induced current is proposed, which can be used to estimate the maximum current value by measuring the displacement value without measuring the current value, thereby realising the same switching signal at the induced current peak. Therefore, we proposed CPCS as:

[CPCS]:  
If  $x(t_0) > 0$ , when  $t = t_0 + \Delta t$ , point Y should be selected.  
If  $x(t_0) < 0$ , when  $t = t_0 + \Delta t$ , point X should be selected.

Here,  $\Delta t$  is the time difference between the current and velocity peaks that can be calculated as follows. Upon substituting Eq. (8) into

Eq. (7) and integrating,  $I$  can be calculated as

$$I = C_1 \exp\left(-\frac{R_0}{L_m^s} t\right) - \frac{\omega b_m U}{\sqrt{R_0^2 + (\omega L_m^s)^2}} \sin\left(\omega t + \gamma + \varphi\right), \quad (11)$$

where

$$\alpha \equiv \frac{R_0}{L_m^s}, \quad \varphi \equiv \tan^{-1}\left(\frac{\alpha}{\omega}\right) = \tan^{-1}\left(\frac{R_0}{\omega L_m^s}\right), \quad (12)$$

where  $C_1$  is constant and  $\varphi$  is the phase difference between the strain and the induced current. Eq. (11) can be differentiated to give

$$\dot{I} = C_2 \exp\left(-\frac{R_0}{L_m^s} t\right) - \frac{\omega^2 b_m U}{\sqrt{R_0^2 + (\omega L_m^s)^2}} \sin\left(\omega t + \gamma + \varphi - \frac{\pi}{2}\right), \quad (13)$$

where  $C_2$  is constant. Because under the first natural vibration mode,  $x$  and  $u$  are proportional and in phase,  $\ddot{u}$  and  $x$  have the same phase relationship. Meanwhile, the sign judgement indices of CPCS and VPCS are opposite. From Eq. (13), we can derive that  $\dot{I}$  and  $x$  have an opposite phase relationship. The phase difference between  $\dot{I}$  and  $u$  is  $\pi/2 - \varphi$ , and  $\Delta t$  can be derived and defined as

$$\Delta t \equiv \frac{\pi}{2\omega} - \frac{1}{\omega} \tan^{-1}\left(\frac{R_0}{\omega L_m^s}\right). \quad (14)$$

By combining VPCS and  $\Delta t$ , CPCS can be realised by measuring the displacement and electrical parameters of the circuit. CPCS can then be rewritten as follows:

[CPCS]:  
If  $x(t_0) > 0$ , when  $t = t_0 + \frac{\pi}{2\omega} - \frac{1}{\omega} \tan^{-1}\left(\frac{R_0}{\omega L_m^s}\right)$ , point Y should be selected.  
If  $x(t_0) < 0$ , when  $t = t_0 + \frac{\pi}{2\omega} - \frac{1}{\omega} \tan^{-1}\left(\frac{R_0}{\omega L_m^s}\right)$ , point X should be selected.

Fig. 6 shows the switching signals obtained using CPCS and VPCS. The switching signals maintain the same frequency as the structural vibration. The CPCS switching signal shows that when the induced current increases, point X is selected; otherwise, point Y is selected. The VPCS switching signal shows that when the displacement is negative, point X is selected; otherwise, point Y is selected.

### 3. Numerical simulation for each harvesting control strategy

To compare the harvesting performances of the control strategies, we simulated the vibration of the cantilevered 10-bay truss structure shown in Fig. 7. Here,  $x$  represents the displacement of the truss structure along the  $X$ -axis. The magnetostrictive transducer was attached to the base of the structure and replaced one bar member in the first bay. The modal damping ratio  $\zeta_i$  of each vibration mode was determined as  $9.7 \times 10^{-3}$ . We analysed this system as a discrete structure and systematically constructed the equation of motion using the finite element method. The parameters of the numerical simulation are shown in Table 1. Each parameter was based on actual experimental setups. The excited vibration mode was the first natural bending vibration mode in the  $X$ -axis, and the excitation point was set at the fourth bay of the truss structure. The motion equation for the magnetostrictive transducer with the truss structure can be derived from Appendix A. All numerical simulations were conducted using MATLAB (MathWorks, Inc., USA).

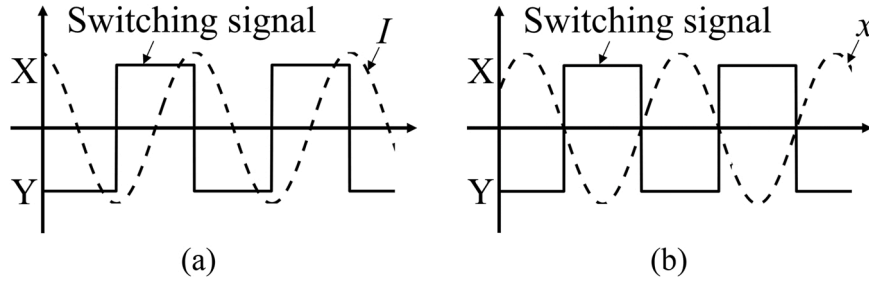


Fig. 6. Switching signals achieved using current peak control strategy (CPCS) and velocity peak control strategy (VPCS).

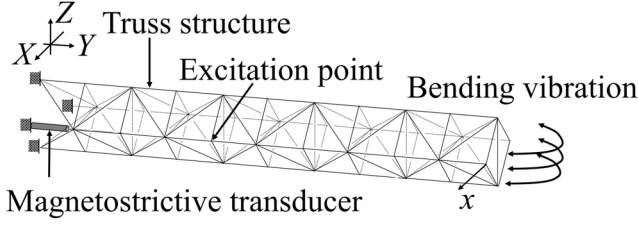


Fig. 7. Truss structure with magnetostrictive transducer.

Table 1  
Parameters of numerical simulation.

Parameter	Value
Stiffness of bar member, $k_t$	$1.99 \times 10^6$ N/m
Ordinary bar member length, $l_o$	$3.80 \times 10^{-1}$ m
Diagonal bar member length, $l_d$	$5.40 \times 10^{-1}$ m
Ordinary bar member mass, $m_o$	$3.57 \times 10^{-2}$ kg
Diagonal bar member mass, $m_d$	$4.63 \times 10^{-2}$ kg
Node mass, $m_n$	$6.79 \times 10^{-2}$ kg
Stiffness at constant current, $k_m^I$	$1.60 \times 10^7$ N/m
Magnetostriction coefficient, $b_m$	$1.15 \times 10^2$ N/A
Inductance at constant velocity, $L_m^S$	$1.90 \times 10^{-1}$ H
Inversion capacitance, $C_i$	$1.00 \times 10^{-6}$ F
Harvesting capacitance, $C_h$	$1.00 \times 10^{-3}$ F
Internal resistance, $R_0$	$2.18 \times 10^1$ $\Omega$
Time step	$1.00 \times 10^{-5}$ s

### 3.1. Harvesting processes and comparison between VPCS and CPCS

We simulated an excitation vibration acting on the truss structure. The motion and circuit equations were calculated using the Runge-Kutta fourth-order method. The initial values  $I_0$  and  $x_0$  were both zero. Because we intended to compare the harvesting performances of VPCS and CPCS, the passive method was performed until its harvested energy reached saturation. The simulation time was 20.0 s, during which the passive method was used for the first 10.0 s and the switching signal from 10.0 s to the end of the experiment. The frequency of the first natural vibration mode was 11.40 Hz.

Fig. 8 shows the time histories of the voltages of  $C_h$  ( $V_h$ ) and  $C_i$  ( $V_i$ ) obtained using CPCS. Within the first 10.0 s of the excitation,  $V_i$  maintained a zero-state because no LC electrical oscillations occurred.  $V_h$  reached a saturation value immediately because the half rectifier passive circuit was conducted. When switching started at 10.0 s,  $V_i$  rapidly increased due to the LC electrical oscillation; in particular, the amplitude peak of  $V_i$  far exceeded the saturation value of  $V_h$  achieved by the passive method in the previous stage. From the time histories of  $V_h$  obtained during switching, we confirmed that the proposed control strategies realised a higher voltage stored in  $C_h$  than the passive method. Fig. 8(b) shows the partial enlargements at 15.0 s, where  $V_i$  forms an oscillation due to the switching. When  $V_i$  is larger than  $V_h$  before the switching, the charge stored in  $C_i$  is transferred to  $C_h$ . Hence, we assume

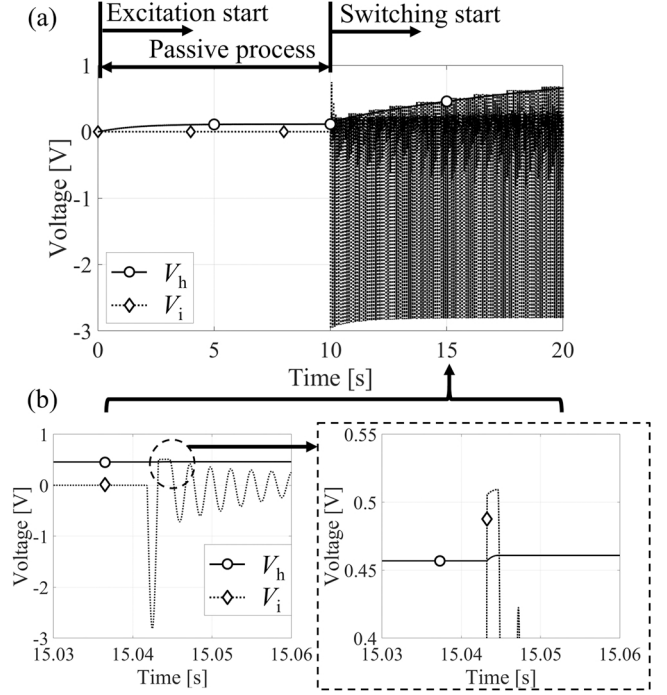


Fig. 8. (a) Time histories of voltages of inversion ( $V_i$ ) and harvesting ( $V_h$ ) capacitors, and (b) their partial enlargements at 15.0 s using current peak control strategy.

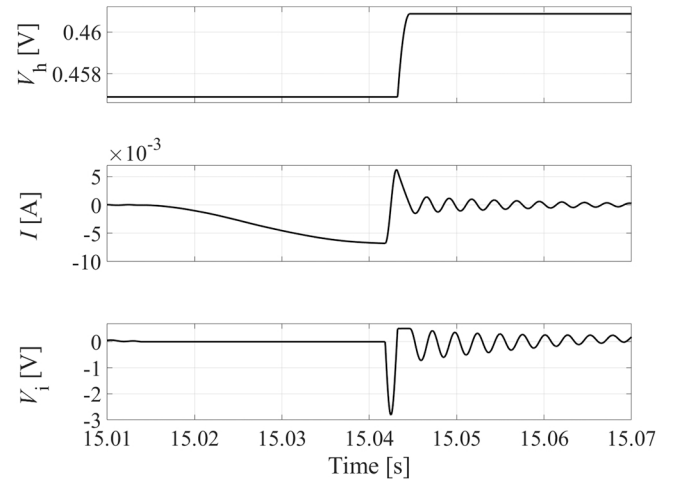
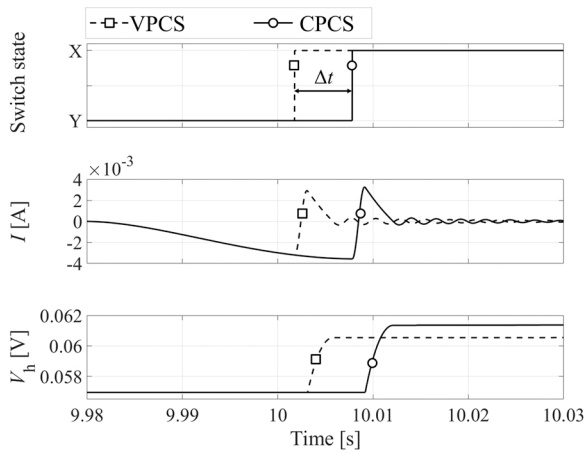


Fig. 9. Time histories of current ( $I$ ) and voltage of inversion capacitor ( $V_i$ ) during switching of connection points (partial enlargement).



**Fig. 10.** Time histories of current ( $I$ ) and voltage of harvesting capacitor ( $V_h$ ) during switching using current peak control strategy (CPCS) and voltage peak control strategy (VPCS) (partial enlargement).

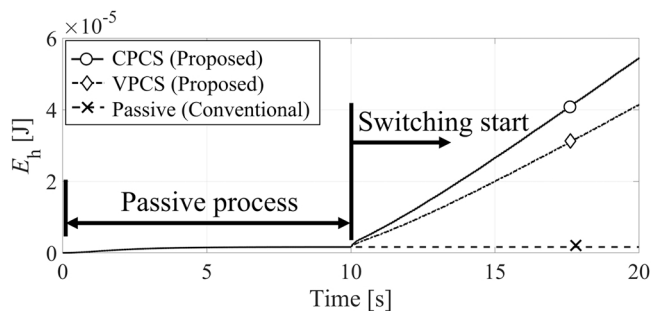
that the amplitude of the LC oscillation determines the final voltage value of  $V_h$ . Because the charge transfer is due to the electric potential difference between two capacitors, the acceleration of  $V_h$  reduces and continuously approaches the amplitude of the LC oscillation.

Fig. 9 shows the partial time histories of  $V_i$ ,  $I$ , and  $V_h$  obtained using both switching control strategies. Because the switching signals changed the connection points,  $I$  inverted from negative to positive due to the LC electrical oscillation, and  $V_i$  reached a maximum value.  $C_h$  collected a certain amount of electrical energy at this stage, and  $V_h$  increased. Next, because of the circuit disconnection, LC electrical oscillation occurred, and  $I$  oscillated and decreased due to  $R_0$  until the magnetostrictive transducer generated a negative current.

Fig. 10 shows the partial time histories of  $V_h$ ,  $I$ , and the switching signals obtained using both switching control strategies. Because of the introduction of  $\Delta t$ , the current peaks of CPCS and VPCS tend to differ. When switching between connection points,  $I$  maintains a higher current state in CPCS than in VPCS. Because of the difference in the magnitude of the initial current at the moment of LC electrical oscillation, the electrical energy stored in the circuit at this time differs. A larger initial current results in a larger amplitude of the electrical energy oscillating in the circuit. Both LC electrical oscillations have the same period that is determined by  $C_i$  and  $L_m^S$ : the larger the amplitude of the LC electrical oscillation, the larger the rate of increase and decrease of  $V_i$ . The switching control strategy with larger increasing and decreasing rates indicates faster charge transport. Therefore,  $V_h$  of CPCS is higher than that of VPCS.

### 3.2. Harvesting performance of each strategy

As mentioned in section 3.1, CPCS can realise a more efficient energy



**Fig. 11.** Time histories of harvested electrical energy  $E_h$  using passive method, current peak control strategy (CPCS), and voltage peak control strategy (VPCS).

**Table 2**

Comparison of harvested energies from numerical simulations.

Harvesting method	Harvested energy [J]
Current peak control strategy (Proposed)	$5.46 \times 10^{-5}$
Voltage peak control strategy (Proposed)	$4.15 \times 10^{-5}$
Passive (Conventional)	$1.63 \times 10^{-6}$

harvesting performance than VPCS. We discuss VPCS, CPCS, and the conventional passive method in this section. To achieve VPCS, we only need the displacement of the 10th bay of the truss structure. Furthermore, by introducing a deliberate time difference based on VPCS, CPCS can be achieved. The harvested electrical energy  $E_h$  is introduced to measure the energy harvesting performance.  $E_h$  is defined as:

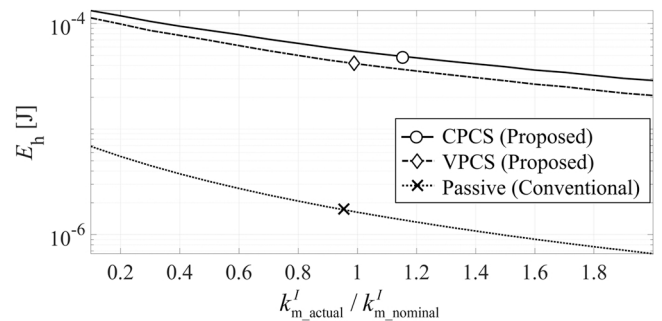
$$E_h \equiv \frac{1}{2} C_h V_h^2. \quad (15)$$

Fig. 11 shows the time histories of  $E_h$  using the passive method, CPCS, VPCS, and the conventional passive method. Over 20.0 s (with switching from 10.0 s), the two switching control strategies realise a significantly higher energy harvesting performance than the conventional passive method. Simultaneously, CPCS has a faster energy increase rate than VPCS, and the energy harvesting performance is also higher. The harvested energy obtained using each switching control strategy is summarised in Table 2. VPCS and CPCS realise 25 and 33 times more energy than the passive method, respectively. Furthermore, the difference achieved in the energy harvesting performance of the two switching control strategies implies that switching timing is a determining factor for energy harvesting performance.

### 3.3. Robustness analysis

Because of changes in material properties due to environmental changes, various unpredictable influences occur in actual harvesting circuits. Therefore, to investigate the robustness of the model error energy harvesting performance, we simulated energy harvesting using intentional model errors based on each harvesting method. For instance, parameters  $k_m^I$  and  $b_m$  are likely to cause differences between values designed under nominal conditions and values obtained in actual systems; therefore, they are valid parameters to manipulate for robustness validation.

Fig. 12 shows  $E_h$  obtained upon varying  $k_m^I$ . Because the exciting force is constant, each switching control strategy realised a higher energy harvesting performance than the passive strategy. Under all conditions, the energy harvesting performance decreases as  $k_m^I$  increases. This indicates that increasing  $k_m^I$  reduces the Villari-Joule effect, thereby reducing the energy harvesting performance of the system. Meanwhile,  $E_h$  decreases as  $k_{m\_actual}^I / k_{m\_nominal}^I$  decreases. Here,  $k_{m\_actual}^I$  and  $k_{m\_nominal}^I$



**Fig. 12.** Harvested electrical energy  $E_h$  using the passive method, current peak control strategy (CPCS), and voltage peak control strategy (VPCS) as a function of errors in the stiffness of the magnetostrictive transducer at constant current,  $k_m^I$ .

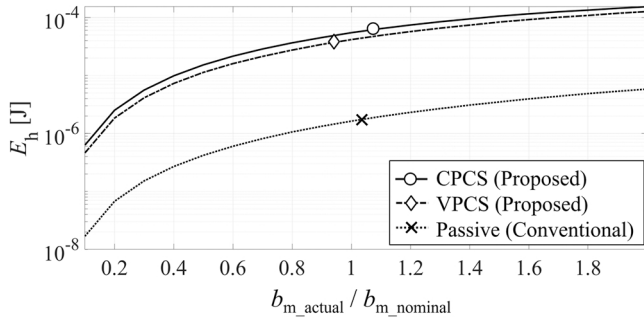


Fig. 13. Harvested electrical energy  $E_h$  using the passive method, current peak control strategy (CPCS), and voltage peak control strategy (VPCS) as a function of errors in the magnetostriction coefficient of the magnetostrictive transducer,  $b_m$ .

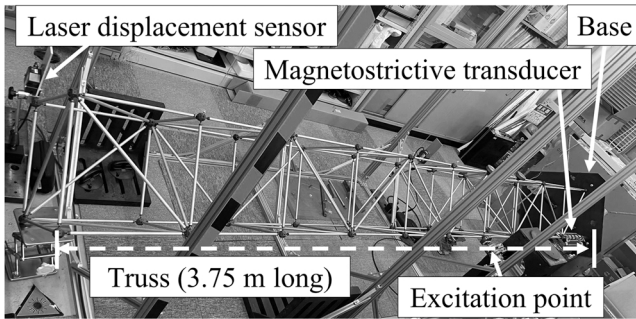


Fig. 14. Magnetostrictive transducer attached in truss structure.

represent the stiffness of the magnetostrictive transducer at a constant current obtained under actual and nominal conditions, respectively. The proposed switching control strategies exhibit similar robustness as the passive method.

Fig. 13 shows  $E_h$  obtained upon varying  $b_m$ . Because the exciting force is constant, each switching control strategy realises higher energy harvesting performance than the passive strategy. Under all conditions, the energy harvesting performance increases as  $b_m$  increases. Moreover,  $E_h$  increases as  $b_{m\_actual}/b_{m\_nominal}$  increases. Here,  $b_{m\_actual}$  and  $b_{m\_nominal}$  represent the magnetostriction coefficient of the magnetostrictive transducer obtained under actual and nominal conditions, respectively.

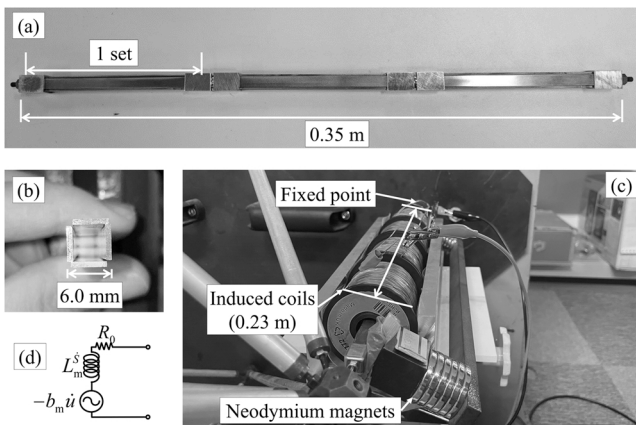


Fig. 15. (a) Top and (b) cross-sectional views of magnetostrictive materials. (c) Magnetostrictive transducer. (d) Electrical model of proposed magnetostrictive transducer:  $L_m^S$ , inductance of magnetostrictive transducer at constant velocity;  $b_m$ , magnetostriction coefficient of magnetostrictive transducer;  $u$ , elongation of the magnetostrictive material;  $R_0$ , internal resistor.

From Eqs. (1) and (2), we observe that increasing  $b_m$  increases the induced voltage, such that the energy harvesting performance is also increased. The proposed switching control strategies exhibit similar robustness as the passive method.

#### 4. Experimental validation of energy harvesting

##### 4.1. Truss structure and equipment setup

To validate the energy harvesting performance of the proposed method in an actual structure, experiments were conducted using a cantilevered 10-bay truss structure with a length of 3.75 m and weight of  $1.32 \times 10^1$  kg. Fig. 14 shows the experimental system. The dimensions of the bar members comprising the truss structure are summarised in Table 1. The bar members were connected by iron nodes. The truss structure was fixed on a base (end support). Two strings were used to counteract gravity such that the truss structure remained straight in the horizontal direction.

The magnetostrictive transducer replaced one ordinary bar member located at the first bay and connected to the fixed end support of the truss structure. Fig. 15 shows the magnetostrictive transducer used. The magnetostrictive material consisted of 12 ready-made Galfenol sheets with 1.0 mm thickness, 6.0 mm width, and  $1.00 \times 10^2$  mm length. Every four sheets formed a square hollow rod with a side length of 6.0 mm, as shown in Fig. 15(b). Three sets of rods comprised the magnetostrictive transducer with  $1.60 \times 10^6$  N/m  $k_m^I$ , as shown in Fig. 15(a). Around the material, a steel beam with permanent magnets contributed to the external magnetic field to improve the Villari-Joule effect. Four cylindrical coils connected in tandem were used to induce a current as the magnetic field was varied, as shown in Fig. 15(c). The entire magnetostrictive transducer weighed 0.6 kg, of which the weight of the material component was 0.1 kg. Additional parameters are shown in Table 1.

Fig. 16 shows a schematic of the experimental setup. The excitation point was at the centre of the truss structure. The laser displacement sensors (LK-030, Keyence Corp.) were used to measure the 10th bay of the truss structure. The laser displacement metre data were used to determine the switching signal. An A/D converter was used to capture the metre data from LK-030 and output the switching signal.

##### 4.2. Experimental process of proposed methods

The process of experimental implementation is depicted in Fig. 17 using a flowchart, and can be described as follows:

- We set up a laser displacement sensor (LK-030, Keyence Corporation) to measure the displacement of the 10th bay node in the  $X$ -axis direction. Using the A/D converter, the analog data from the laser displacement sensor were converted to digital data.
- The timing of the displacement extremum can be detected by the PC programme. Here, according to the proposed control strategies,

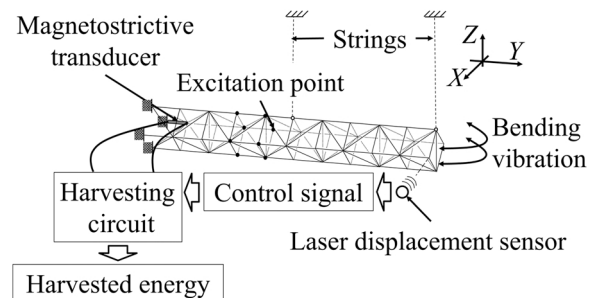


Fig. 16. Experimental setup of truss structure with magnetostrictive transducer.



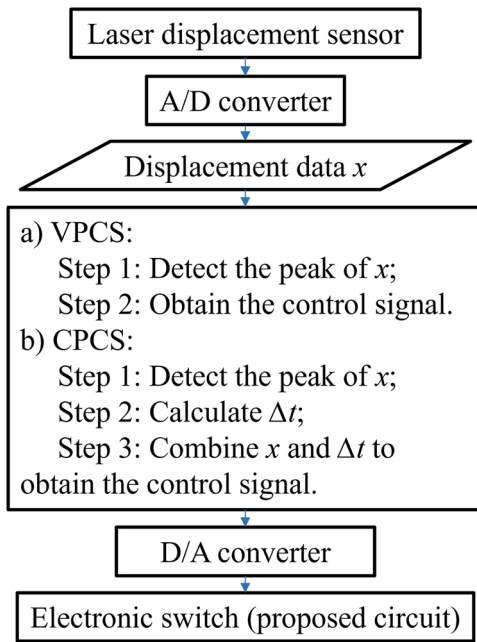


Fig. 17. Experimental implementation process.

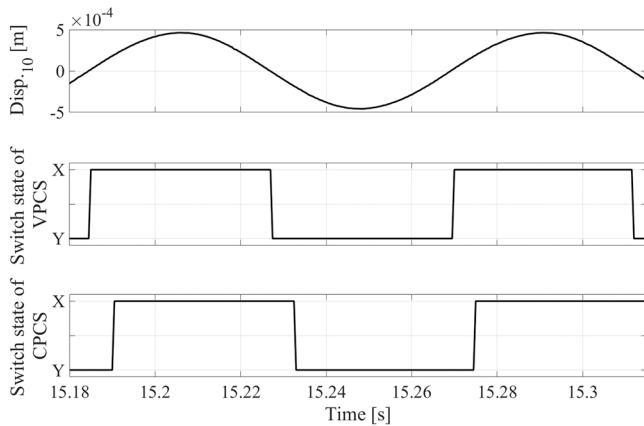


Fig. 18. Time histories of displacement of 10th bay and switching signal for current peak control strategy (CPCS) and voltage peak control strategy (VPCS).

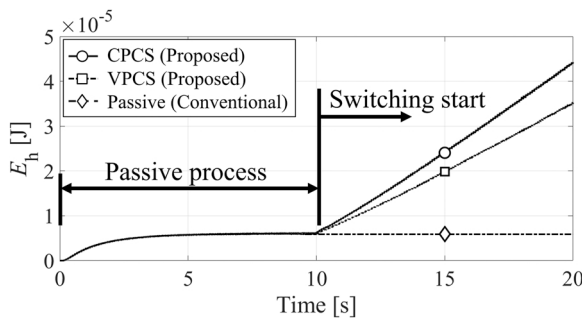


Fig. 19. Time histories of harvested energy ( $E_h$ ) for passive method, current peak control strategy (CPCS), and voltage peak control strategy (VPCS) obtained from experiments.

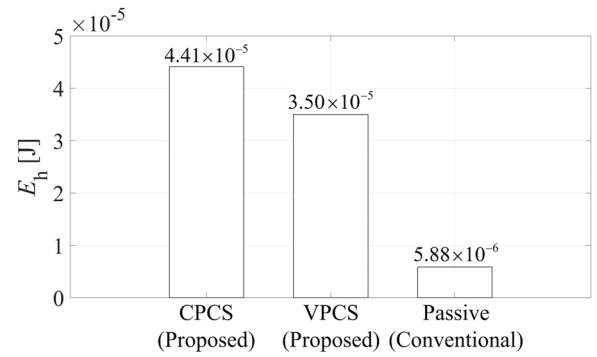


Fig. 20. Harvested energy ( $E_h$ ) for passive method, current peak control strategy (CPCS), and voltage peak control strategy (VPCS) over 20.0 s obtained from experiments.

displacement extremum and offering  $\Delta t$ , which was calculated from known parameters of the circuit elements.

- c) Using the PC programme, we selected the connection of the electronic switch at each moment. These connection data were expressed as control signals.
- d) The control signal was converted into a pulse voltage by the A/D converter and output to the electronic switch.
- e) The voltage of the harvesting capacitor was measured to evaluate the harvesting performance.

### 4.3. Energy harvesting performance comparison

In the experiments, the excitation direction was along the X-axis, as shown in Fig. 16. The truss structure was excited in the bending first natural vibration mode (11.81 Hz). For comparison, the same experiments were conducted using the conventional passive method for 10.0 s and then conducted with VPCS and CPCS for 10.0 s

Fig. 18 shows the time histories of the displacement of the 10th bay and the switching signal obtained using both switching control strategies. The switching frequency was the same as the structural vibration frequency. Meanwhile, compared with the VPCS signal, a fixed phase difference shifted the CPCS signal.

Fig. 19 shows the time histories of the harvested energy obtained using both switching control strategies. As mentioned in section 2.1, the kinetic energy from structural vibrations is converted into electrical energy. The time histories of  $E_h$  show that CPCS and VPCS harvest more electrical energy than the passive method and CPCS realises a higher harvesting performance than VPCS.

Lastly, the harvested energy obtained using both switching control strategies is summarised in Fig. 20. We found VPCS and CPCS realise 5.9 and 7.5 times more energy than the passive method, respectively.

## 5. Conclusion

In this study, two new switching control strategies for energy harvesting that combine a magnetostrictive transducer and electric circuits were proposed. The proposed circuit includes two capacitors, two diodes, and an electronic switch. As a result of the instantaneous switching between points, LC electrical oscillation occurs over a short period. Simultaneously, this short-term LC oscillation increases the voltage across the inverse capacitor to a maximum value, which promotes the continuous transfer of charge to the harvesting capacitor. The proposed method utilises the instantaneous extreme voltages generated by LC electrical oscillations to achieve large-scale and efficient energy harvesting performances. Approximately 45  $\mu$ J of energy were harvested over 20 s.

Through numerical simulations, the feasibility and robustness of each proposed switching control strategy were predicted. To validate

displacement data were used to obtain the control signal through the PC programme. We performed VPCS by detecting the displacement extremum. Similarly, we performed CPCS by detecting the

their energy harvesting performance, experiments were performed using a cantilevered truss structure. The results showed that the proposed strategies were more effective than the passive method (no switching) with similar robustness. In the experimental environment, CPCS harvested 7.5 times more energy than the conventional passive method. The amount of harvested energy observed in this study was relatively small for use in practical applications, and the switching control strategies can only be used for single-mode vibrations. Nevertheless, they achieved a substantial increase in energy harvesting compared to the passive method without needing to alter the properties of the raw materials involved in the system.

In future work, the switching control strategies will be extended to multimodal vibration conditions. Additionally, a new design for a more efficient and lightweight transducer will be investigated to satisfy the energy supply requirement of the electronic switches and space-constrained application environment.

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## Appendix A

In this study, the magnetostrictive transducer was attached in the infrastructure truss structure, as shown in Fig. 7. During horizontal bending vibration, the deformation amplitude at the fixed end of the truss is larger, thus achieving more effective energy harvesting. Compared with the dimensions of the truss structure, the magnetostrictive transducer is a relatively small-size energy converter. When the local dynamics of a magnetostrictive transducer are negligible, the position coordinate vector  $\mathbf{x}$  of each node in the global coordinate system is related to the displacement vector  $\mathbf{u}_i$  in the  $i$ -th local coordinate system via a transformation matrix  $\mathbf{H}_i$ :

$$\mathbf{u}_i = \mathbf{H}_i \mathbf{x}, \quad (\text{A.1})$$

where  $\mathbf{u}_i$  is defined as  $[u_k, u_j]^T$ , and  $u_k$  and  $u_j$  represent the positions of the  $k$ -th and  $j$ -th nodes connected to the  $i$ -th bar member, respectively. The axial force  $P_i$  applied to the  $i$ -th bar member can be described as

$$P_i = k_{ti}(u_k - u_j), \quad (\text{A.2})$$

where  $k_{ti}$  represents the stiffness of the  $i$ -th bar member. The potential energy of a bar member can be described as

$$W_{mi} = (\mathbf{G}\mathbf{H}_i \mathbf{x} - b_{ci} I_i) \frac{k_{ci}}{2} (\mathbf{G}\mathbf{H}_i \mathbf{x} - b_{ci} I_i), \quad (\text{A.3})$$

where

$$k_{ci} \equiv \frac{k_{mi}^J k_{ti}}{k_{mi}^J + k_{ti}}, \quad b_{ci} \equiv \frac{b_{mi}}{k_{mi}^J}, \quad \mathbf{G} = [1 \quad -1]. \quad (\text{A.4})$$

Here,  $I_i$ ,  $b_{mi}$ , and  $k_{mi}^J$  represent the current, magnetostriction coefficient, and constant-current stiffness of the  $i$ -th magnetostrictive material, respectively. The total potential energy  $W_{\text{total}}$  of the truss structure can be expressed as

$$W_{\text{total}} = \sum_{i=1}^{N_m} W_{mi} + \sum_{i=N_m+1}^{N_t} W_{ti} = \sum_{i=1}^{N_m} (\mathbf{G}\mathbf{H}_i \mathbf{x} - b_{ci} I_i) \frac{k_{ci}}{2} (\mathbf{G}\mathbf{H}_i \mathbf{x} - b_{ci} I_i) + \sum_{i=N_m+1}^{N_t} (\mathbf{G}\mathbf{H}_i \mathbf{x}) \frac{k_{ti}}{2} (\mathbf{G}\mathbf{H}_i \mathbf{x}). \quad (\text{A.5})$$

We suppose that  $m_l$  is the mass concentrated at the  $l$ -th node. The total kinetic energy  $T_{\text{total}}$  can then be described as

$$T_{\text{total}} = \sum_{l=1}^{N_{\text{node}}} T_l = \frac{1}{2} \dot{\mathbf{x}}^T \mathbf{M} \dot{\mathbf{x}}, \quad (\text{A.6})$$

where the mass matrix of structure  $\mathbf{M}$  is defined as

$$\mathbf{M} \equiv \text{blockdiag}[\mathbf{M}_1, \mathbf{M}_2, \dots, \mathbf{M}_{N_{\text{node}}}], \quad (\text{A.7})$$

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## CRedit authorship contribution statement

**An Li:** Conceptualization, Methodology, Validation, Writing – original draft, Funding acquisition. **Keiju Goto:** Methodology, Writing – review & editing. **Yuusuke Kobayashi:** Methodology, Writing – review & editing. **Yushin Hara:** Investigation, Writing – review & editing. **Yu Jia:** Writing – review & editing. **Yu Shi:** Writing – review & editing. **Constantinos Soutis:** Writing – review & editing. **Hiroki Kurita:** Writing – review & editing. **Fumio Narita:** Writing – review & editing, Funding acquisition. **Keisuke Otsuka:** Investigation, Writing – review & editing. **Kanjiro Makihara:** Supervision, Funding acquisition.

## Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Data Availability

Data will be made available on request.

$$\mathbf{M}_l \equiv \text{diag}[m_l, m_l, m_l]. \quad (\text{A.8})$$

From Eqs. (A.5) and (A.6), the motion equation for the magnetostrictive transducer can be derived using Hamilton's principle as

$$\mathbf{M}\ddot{\mathbf{x}} + \mathbf{D}\dot{\mathbf{x}} + \mathbf{K}\mathbf{x} = \mathbf{B}_m\mathbf{I} + \mathbf{f}_{\text{ext}}, \quad (\text{A.9})$$

where  $\mathbf{x}$  represents the position coordinates of each node;  $\mathbf{f}_{\text{ext}}$  is the multiple external force matrix; and  $\mathbf{M}$ ,  $\mathbf{D}$ ,  $\mathbf{K}$ ,  $\mathbf{B}_m$ , and  $\mathbf{I}$  are the mass, damping coefficient, stiffness at constant current, magnetostriction-stiffness composite coefficient, and induced current of magnetostrictive transducer matrices, respectively.

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